## On the nature of spacetime singularities

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### The singularity theorem by Roger Penrose (1965)

- during the first five decades of Einstein's theory of gravity, singular behavior popped up in many of the physically relevant solutions
- singularity theorems of Penrose and Hawking
  - they predict, under a wide range of physically plausible conditions, the existence of **spacetime singularities**, more precisely, they prove that
  - generic spacetimes describing the gravitational collapse of stars (1965) or the expanding universe (1966) are **causal geodesically incomplete**

• **singularity**  $\iff$  **incompleteness**: they are used as synonyms

- in the case of the well-known solutions, the incomplete causal geodesics do indeed terminate on "curvature singularities" (e.g.,  $R_{abcd}R^{abcd}$  ...)
- BUT, do we know for sure that the anticipated curvature blow up or any other violent behavior will always occur?
- NO, NOT REALLY !!! NOT YET !!!
- $\bullet \ \exists$  inconsistency in the use of notions: incompleteness and singularities
- outline an argument: a step towards the reduction of this discrepancy

## Spacetime as a Cauchy development?

#### Penrose's strong cosmic censorship conjecture

- $(M, g_{ab})$ : M is a smooth paracompact, connected, orientable manifold endowed with a  $C^{\infty}$  metric  $g_{ab}$  of Lorentzian signature and time orientation
  - $C^\infty$  only for mathematical conveniences...
  - what is the most suitable differentiability class to be used? (return later)
- tacitly a spacetime is always assumed to represent all the events compatible with the history of the investigated physical system
- a celebrated result by Choquet-Bruhat and Geroch (1969) guarantees (assuming smoothness) the existence of a maximal Cauchy development
- $\bullet$  the situation is more complex:  $\exists$  causal geodesically incomplete spacetimes:
  - they contain as a part "the maximal" Cauchy development, and

• they can be continued beyond the Cauchy horizon: no curvature blowup occurs (think of max. anal. extensions of Kerr, Taub-NUT spacetimes where the ability to predict the future of data given on an inextendible initial data surface, breaks down)

• **Penrose's strong cosmic censorship conjecture:** the maximal Cauchy development of a **generic** compact or asymptotically flat initial data is never part of a larger spacetime. (if true, *A* Cauchy Horizon, no extension beyond)

## How could we show that something violent is to happen?

## • What do we have by hand?

- singularity theorems by Penrose and Hawking
- the existence of maximal Cauchy development Choquet-Bruhat and Geroch
- strong cosmic censorship conjecture by Penrose

#### • What about the following argument by contradiction?

- consider a causal geodesically incomplete spacetime that is the maximal Cauchy development of a **generic** compact or asymptotically flat initial data, i.e. **it is inextendible**
- assume that nothing violent happens while approaching the "ideal endpoint" of any of the incomplete causal geodesics
- **show** that the considered maximal Cauchy development **can be extended**; *leading to the contradiction:* **it is not maximal, it is extendible**
- in turn, we get then that causal geodesic incompleteness of the maximal Cauchy development must be accompanied by a singular behaviour
- !!! ...to show that the original spacetime is part of a larger one... we also need some results on *spacetime extensions*: CJC Clarke mid 70'; IR 1993, 2010

## How do we extend a spacetime? manifold & metric

- Definition: Consider (M, g<sub>ab</sub>) and (M̂, ĝ<sub>ab</sub>) the differential structure of which are at least of class C<sup>X</sup> (C<sup>∞</sup>?), respectively. (X not specified yet!) A map Φ : (M, g<sub>ab</sub>) → (M̂, ĝ<sub>ab</sub>) is said to be a C<sup>X</sup>-isometric embedding if Φ is a C<sup>X</sup>-diffeomorphism between M and Φ[M] ⊂ M̂ such that it carries the metric g<sub>ab</sub> into ĝ<sub>ab</sub>|<sub>Φ[M]</sub>, i.e. Φ<sup>\*</sup>g<sub>ab</sub> = ĝ<sub>ab</sub>|<sub>Φ[M]</sub>.
- Definition:  $(\widehat{M}, \widehat{g}_{ab})$  is called to be a  $\mathbf{C}^{\mathbf{X}}$ -extension of  $(M, g_{ab})$  if  $\Phi$  is an isometric embedding such that  $\Phi[M]$  is a proper subset of  $\widehat{M}$



# What is the most suitable differentiability class?

- ${\, \bullet \,}$  smooth or even  $C^2$  may be too much to be required
- general relativity is a physical theory: the wider the class of metrics allowed the wider will be the class of physical processes that can be investigated; ??? field equations ???
- Geroch and Traschen: (1987) the widest possible class of metrics such that the Riemann, Einstein and Weyl tensors are well-defined as distributions
  - the space of regular metrics:
    - $g_{ab}$  locally bounded
    - with locally bounded inverse  $g^{ab}$
    - $\bullet$  the weak first derivatives " $\partial_c g_{ab}$ " are locally square-integrable
  - C<sup>0</sup> regular metrics:
    - if a regular metric  $g_{ab}$  is also **continuous** it can be approximated by sequences of smooth metrics  $\{^{(i)}g_{ab}\}$  such that the corresponding curvature tensors  $\{^{(i)}R_{abc}{}^d\}$  do converge in  $L_2$ -norm to the curvature distribution  $R_{abc}{}^d$  of the continuous regular metric  $g_{ab}$

# The $C^0$ G-T regular metrics are still too rough

- Some of the basic concepts should be well-defined: the class of locally Lipschitz,  $C_{loc}^{0,1} (= C_{loc}^{1-})$ , metrics is distinguished
- C<sup>0,1</sup><sub>loc</sub> differentiability suffices for many key results of the C<sup>∞</sup> causality theory
   global hyperbolicity makes sense
   Chrusciel-Grant (2012), Sämann (2016)
- **global** existence and uniqueness to linear field equations with sources and with  $g_{ab}$  of  $C_{loc}^{0,1}$  (Chrusciel-Grant (2012), Sanches-Vickers 2017)
  - ∃ maximal Cauchy development for the vacuum Einstein equations requires metric with critical Sobolev exponent s = 2 (Klainerman-Rodnianski 2001)
- $C_{loc}^{0,1}$  is the weakest condition under which "pointwise" differential geometry is possible (CJS Clark 1982)
  - $C_{loc}^{0,1}$ -geodesic as uniform accumulation of sequences of smooth  ${}^{(i)}g_{ab}$ -geodesics
  - weak solutions to Jacobi's equation make sense (a.e.) along timelike geodesics
- physically interesting solutions with  $C_{loc}^{0,1}$  regular metrics:
  - gravitational shock waves
  - thin mass shells
  - solutions with pressure free matter with geodesic flow lines having two- or three-dimensional caustics

## What type of result are we looking for?

- **Theorem:** Consider a timelike geodesically incomplete spacetime  $(M, g_{ab})$  that is the maximal globally hyperbolic development of a generic compact or asymptotically flat initial data and that is inextendible as a Lorentzian manifold within the class of  $C_{loc}^{0,1}$  metrics. Then, there exists an incomplete timelike geodesic that terminates on a **parallelly propagated curvature singularity**.
- in what follows the metric will be assumed to be smooth, with the understanding that  $C^{0,1}_{loc}$  metrics can be approximated by sequences of  $C^{\infty}$  metrics
- the papers below provide some hints on results relevant for the smooth setup
  - Rácz, I. (1993): *Spacetime extensions I.*, Journal of Mathematical Physics **34**, 2448-2464
  - Rácz, I. (2010): Space-time extensions II, Classical and Quantum Gravity 27, 155007, Selected by the Editorial Board of Classical and Quantum Gravity as part of the journal's Highlights of 2010, in 2011.

### 1st step:

- assume that  $(M, g_{ab})$  is a timelike geodesically incomplete, inextendible spacetime and that the **tidal force tensor components are bounded**
- let  $\gamma: (t_1, t_*) \to M$  be one of the future incomplete timelike geodesics
  - $\bullet\,$  choose  ${\cal U}$  be a neighbourhood of a final segment of  $\gamma\,$
  - the aim is to find  $\hat{\mathcal{U}}$  and  $\phi : \mathcal{U} \to \hat{\mathcal{U}}$  such that  $\hat{\mathcal{U}}$  is comprised by the union of  $\phi[\mathcal{U}]$  and "a neighbourhood of the endpoint of  $\phi \circ \gamma$ " (not yet constructed !!!)



## 2nd step:



(M̂, ĝ<sub>ab</sub>): defined by gluing (M, g<sub>ab</sub>) and (Ũ, g̃<sub>ab</sub>) at their "common parts"
M̂ is the quotient space M̂ = (M ∪ Ũ)/φ, where x ∈ M is equivalent to y ∈ Ũ if φ(x) = y; M̂ is Hausdorff with respect to the quotient topology
the metric ĝ<sub>ab</sub> on M̂ is determined by g<sub>ab</sub> and g̃<sub>ab</sub>

## The construction of $\mathcal{U}$

- to apply Whitney's extension results: cover  ${\mathcal U}$  by a single coordinate patch
  - let  $\gamma : (t_1, t_*) \to M$  be an incomplete timelike geodesic &  $p = \gamma(t_0)$ : t is the proper time parameter along  $\gamma$ , with tangent  $v^a = (\partial/\partial t)^a$ ; (t is an affine par.)
  - choose  $\Sigma$  to be the hypersurface generated by spacelike geodesics starting at  $p=\gamma(t_0)$  with tangent orthogonal to  $v^a$



- choose an orthonormal frame  $\{e^a_{(\mathfrak{a})}\}$ ,  $e^a_{(4)} = v^a$  at p and extend it onto  $\Sigma$
- consider the 3-par. family of timelike geodesics starting at  $\Sigma$  with  $v^a = e^a_{(A)} \dots$
- denote by  $\Gamma$  the "synchronized" 3-par. family of timelike geodesics
- choose  $\mathcal{U}$  to be ruled by members of  $\Gamma$ : extend  $\{e^a_{(\mathfrak{a})}\}$  onto  $\mathcal{U}$  by par.prop...

## Can $\Sigma$ be chosen such that $\mathcal{U}$ covers the final segments?

• whenever the tidal force components

$$R_{abcd} e^a_{(\mathfrak{a})} v^b e^c_{(\mathfrak{b})} v^d \qquad \qquad (\text{where } v^a = e^a_{(4)})$$

of the Riemann tensor are bounded in synchronized, parallelly propagated orthonormal frames  $\{e^a_{(\mathfrak{a})}\}$ , along the synchronized 3-parameter congruence  $\Gamma$  (we tacitly assumed that no topological obstruction [alg. special cases] occur)

 Σ can be chosen such that U is ruled by final segments of members of the timelike geodesic congruence Γ, and such that no conjugate point to Σ occurs



 $\bullet$  then Gaussian coordinates  $(x^1,x^2,x^3,x^4=t)$  can be defined on  ${\mathcal U}$ 

# How to extend the metric from $\phi[\mathcal{U}] \subset \mathbb{R}^4$ ?

#### • How does one extend a function?

- Consider a real-valued  $C^m$ -function  $\mathcal{F}$  defined on a bounded subset  $\mathscr{A}$  of  $\mathbb{R}^n$ .
- How can we tell whether there exists  $\widetilde{\mathcal{F}} \in C^m(\mathbb{R}^n)$  such that  $\widetilde{\mathcal{F}} = \mathcal{F}$  on  $\mathscr{A}$ ? [studied by Hassler Whitney ~1930 (see also C Fefferman (2005))]
- **Definition:** a point set  $\mathscr{A} \subset \mathbb{R}^n$  is said to possess **the property**  $\mathscr{P}$  if there is a positive real number  $\omega$  such that for any two points x and y of  $\mathscr{A}$  can be joined by a curve in  $\mathscr{A}$  of length  $L \leq \omega \cdot \rho(x, y)$ , where  $\rho(x, y)$  denotes the Euclidean distance of the points  $x, y \in \mathbb{R}^n$ .
- **Theorem:** [adopting Whitney's results (1934)] Assume that  $\mathscr{A} \subset \mathbb{R}^n$  is bounded and it possesses property  $\mathscr{P}$ , and let  $\mathcal{F}(x^1, ..., x^n)$  be of class  $C_{loc}^{0,1}$ in  $\mathscr{A}$ . Suppose that  $\mathcal{F}(x^1, ..., x^n)$  can be defined on the boundary  $\partial \mathscr{A}$  of  $\mathscr{A}$ such that it is of class  $C_{loc}^{0,1}$  on  $\overline{\mathscr{A}} = \mathscr{A} \cup \partial \mathscr{A}$ . Then there exists an extension  $\widetilde{\mathcal{F}}$  of  $\mathcal{F}$  onto  $\mathbb{R}^n$  such that  $\widetilde{\mathcal{F}}$  is (at least) of class  $C_{loc}^{0,1}$ .
- as the spacetime is a maximal globally hyperbolic development  $\Rightarrow \phi[\mathcal{U}] \subset \mathbb{R}^4$  possesses property  $\mathscr{P}$

## We still need to extend the metric from $\phi[\mathcal{U}]$ !!!

 $\bullet$  in the Gaussian coordinates  $(x^1,x^2,x^3,x^4=t)$  on  $\mathcal{U}:$  the metric reads as

$$ds^2 = -dt^2 + g_{ij} \, dx^i dx^j$$

• Whitney's theorem  $\Rightarrow$  the extendibility of functions defined on bounded subsets of  $\mathbb{R}^4 \Rightarrow$  the components of the metric  $g_{ab}$  can be extended such that its extension  $\tilde{g}_{\alpha\beta}$  is of class  $C^{0,1}_{loc}$  if components  $g_{ij}$  can be shown to be locally Lipschitz functions on the closure  $\overline{\phi[\mathcal{U}]}$  of  $\phi[\mathcal{U}]$ , which holds

• 
$$\Rightarrow$$
 if the "t-derivatives" of  $g_{ij} = g_{ab} E^a_{(i)} E^b_{(j)}$ , where  $E^a_{(i)} := (\partial/\partial x^i)^a$ ,

$$\partial_t g_{ij} = v^e \nabla_e [g_{ab} \, E^a_{(i)} E^b_{(j)}] = g_{ab} \left[ \left( v^e \nabla_e E^a_{(i)} \right) E^b_{(j)} + E^a_{(i)} \left( v^e \nabla_e E^b_{(j)} \right) \right]$$
 are uniformly bounded along the members of  $\Gamma$ 

- it suffices to show that the norms  $||E^a_{(\alpha)}||$  and  $||v^e \nabla_e E^a_{(\alpha)}||$  of the coordinate basis fields  $E^a_{(i)}$ , and also that of  $v^e \nabla_e E^a_{(i)}$  are uniformly bounded on  $\phi[\mathcal{U}]$ 
  - where the norm  $||X^a||$  of a vector field  $X^a$ , with respect to a synchronized orthonormal basis field  $\{e^a_{(a)}\}$  and a Lorentzian metric  $g_{ab}$ , is defined as

$$\|X^a\| = \sqrt{\sum_{b=1}^4 \left[g_{ab}X^a e^b_{(b)}\right]^2}$$

• notably, the boundedness can be guaranteed using the Jacobi equation

$$v^e \nabla_e \left( v^f \nabla_f E^a_{(\alpha)} \right) = R_{efg}{}^a v^e E^f_{(\alpha)} v^g$$

and applying our (indirect) assumption guaranteeing the boundedness of the tidal force components of the Riemann tensor, along the members of  $\Gamma$ 

- combining all the above partial results  $\Rightarrow$  **Theorem:** Consider a timelike geodesically incomplete spacetime  $(M, g_{ab})$ that is the maximal globally hyperbolic development of a generic compact or asymptotically flat initial data and that is inextendible within the class of  $C_{loc}^{0,1}$  spacetimes. Then, there exists an incomplete timelike geodesic that terminates on a **parallelly propagated curvature singularity**.
- $\Rightarrow$  at least one of the tidal force tensor components of the Riemann tensor blows up

### Thanks for your attention