

# Hyperboloidal massless scalar field in 3D

Alex Vano-Vinuales



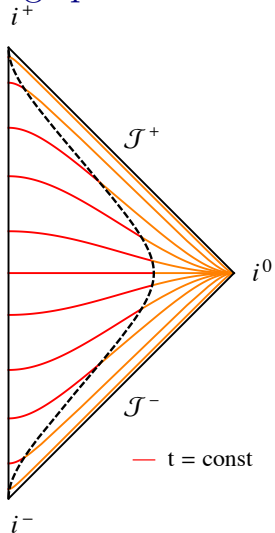
CENTRA, Instituto Superior Técnico



Spanish-Portuguese Relativity Meeting (EREP), online (Aveiro) -

16th September 2021

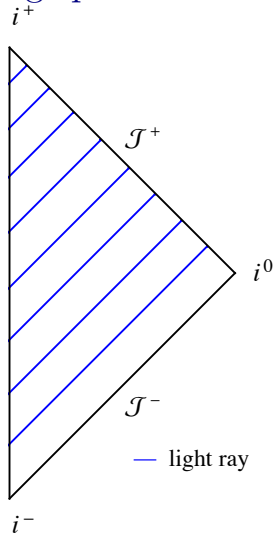
# Slicing spacetime to reach future lightlike infinity



Standard slicing options for the **initial value formulation** of the Einstein equations, to solve them as an evolution in time:

- Standard Cauchy slices

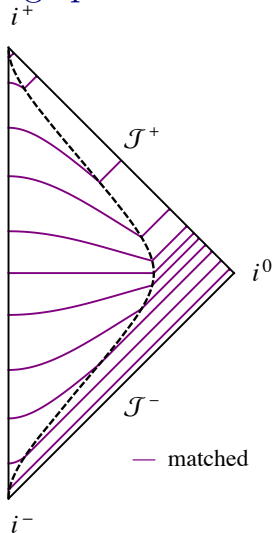
# Slicing spacetime to reach future lightlike infinity



Standard slicing options for the **initial value formulation** of the Einstein equations, to solve them as an evolution in time:

- Standard Cauchy slices
- Null slices

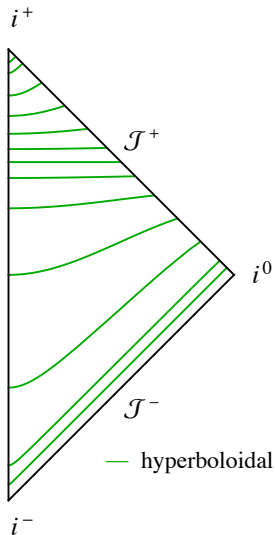
# Slicing spacetime to reach future lightlike infinity



Standard slicing options for the **initial value formulation** of the Einstein equations, to solve them as an evolution in time:

- Standard Cauchy slices
- Null slices
- Cauchy-Characteristic matching / extraction

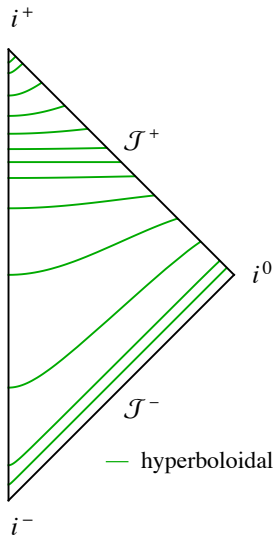
## Slicing spacetime to reach future lightlike infinity



Standard slicing options for the **initial value formulation** of the Einstein equations, to solve them as an evolution in time:

- Standard Cauchy slices
- Null slices
- Cauchy-Characteristic matching / extraction
- Hyperboloidal slices

## Slicing spacetime to reach future lightlike infinity



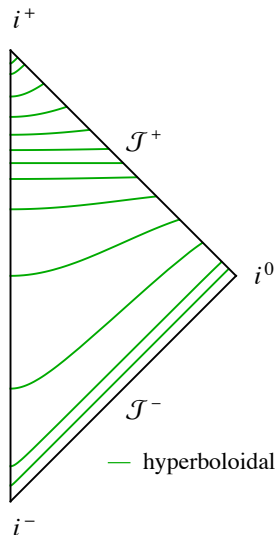
Standard slicing options for the **initial value formulation** of the Einstein equations, to solve them as an evolution in time:

- Standard Cauchy slices
- Null slices
- Cauchy-Characteristic matching / extraction
- Hyperboloidal slices

Advantages of the hyperboloidal approach:

- Extraction at  $\mathcal{J}^+$ , no approximations.
- Slices **spacelike** & **smooth** everywhere.
- More **resolution** for the central part.

# Slicing spacetime to reach future lightlike infinity



**Future null infinity** ( $\mathcal{I}^+$ ) is a region of spacetime of interest

- for the study of **global properties** of spacetimes and
- for the extraction of **gravitational waves** (only well described at  $\mathcal{I}^+$ , where **observers** are located).

A possible approach: Penrose's **conformal compactification** of spacetime.

The physical metric  $\tilde{g}_{\mu\nu}$  is rescaled

$$g_{\mu\nu} \equiv \Omega^2 \tilde{g}_{\mu\nu}, \quad (1)$$

with  $\Omega|_{\mathcal{I}^+} = 0$  to keep  $g_{\mu\nu}$  finite there.

## Equations for the conformal metric

The Einstein equations including cosmological constant are

$$G[\tilde{g}]_{ab} + \tilde{g}_{ab}\Lambda = 8\pi T[\tilde{g}]_{ab}.$$

Expressing them in terms of the rescaled metric  $g_{ab} = \Omega^2 \tilde{g}_{ab}$  gives

$$G[g]_{ab} + \frac{2}{\Omega}(\nabla_a \nabla_b \Omega - g_{ab} \square \Omega) + \frac{3}{\Omega^2} g_{ab} (\nabla_c \Omega)(\nabla^c \Omega) + \frac{1}{\Omega^2} g_{ab} \Lambda = 8\pi T[\frac{g}{\Omega^2}]_{ab}.$$

Extra **formally divergent terms** at  $\mathcal{I}$  appear in the equations.

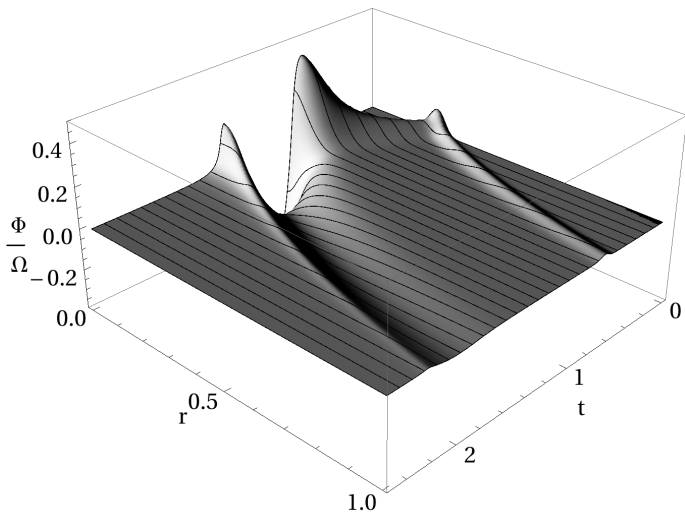
Evaluating the previous expression multiplied by  $\Omega^2$  at  $\mathcal{I}$  yields

$$(\nabla_c \Omega)(\nabla^c \Omega)|_{\mathcal{I}} = -\Lambda, \quad \text{so that}$$

- $\Lambda = 0$ :  $\mathcal{I}$  is **null**,
- $\Lambda > 0$ :  $\mathcal{I}$  is **spacelike**,
- $\Lambda < 0$ :  $\mathcal{I}$  is **timelike**.



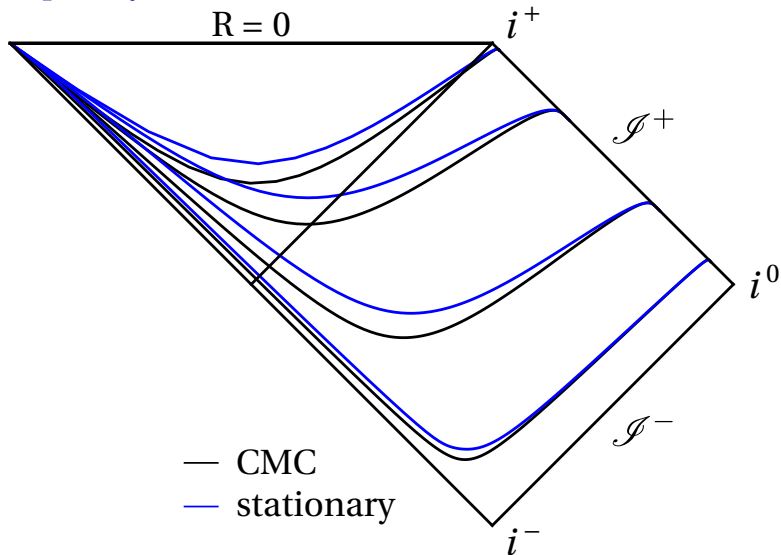
# Scalar field



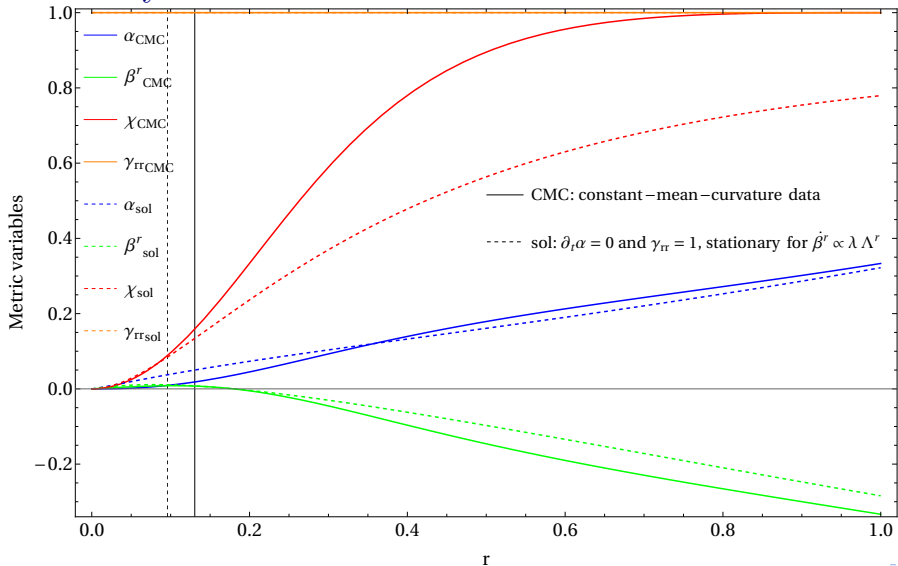
Vañó-Viñuales, Husa and Hilditch. *Class. Quant. Grav.* 32.17 (2015).

Evolution:  $\chi$ ,  $\tilde{K}$ ,  $\alpha$ ,  $\beta^r$ ,  $\Phi/\Omega$  

# Trumpet dynamics



# Stationary initial data



## Wave equation as toy model

Einstein equations are complicated and non-linear, start with simpler models:

### This approach

- Non-conformally invariant:

$$\tilde{\square}\tilde{\Phi} = 0 \quad \rightarrow$$

$$\square\Phi + \Phi \left( \frac{\square\Omega}{\Omega^2} - \frac{2\nabla_a\Omega\nabla^a\Omega}{\Omega} \right) = 0$$

- 2nd order in space
- no matching on slice
- only Minkowski background, Schwarzschild trumpet planned
- by default use spherical coords

### LlamaWaveHyperboloidal\*

- Conformally invariant ( $\tilde{\Phi} = \Omega\Phi$ ):

$$(\tilde{\square} - \tilde{R}/6)\tilde{\Phi} = 0 \quad \rightarrow$$

$$\left( \square - \frac{R}{6} \right) \Phi = 0$$

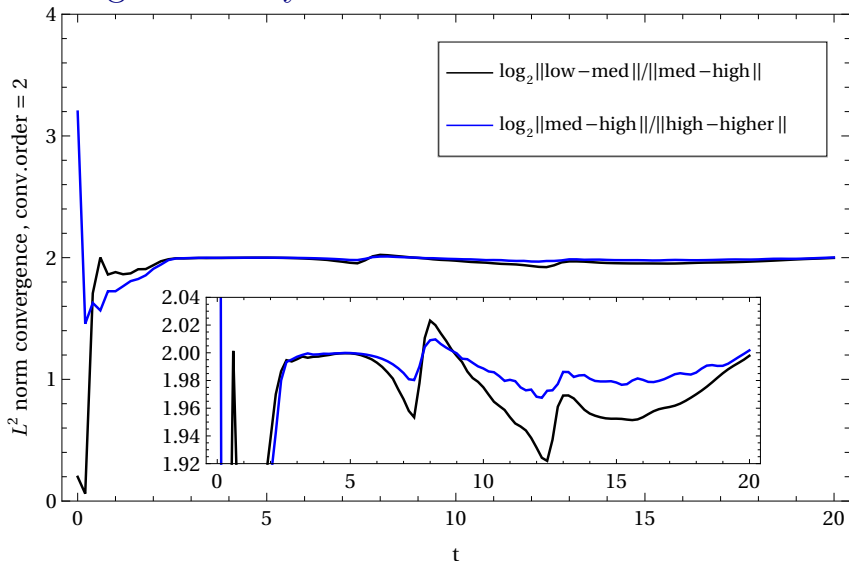
- 1st order, flux-conservative
- matching: inner Cauchy slice and outer hyperboloidal one
- up to Kerr background
- Cartesian coords matched

\*Thorn for the Einstein Toolkit code: [Jasiulek](#). *Class. Quant. Grav.* 29 (2012).

# Massless scalar field equation

Evolution of the massless wave equation on a hyperboloidal slice of Minkowski spacetime using curvilinear coordinates in NRPy+ **l** **r**.

# Convergence: axisymmetric initial data and evolution



# Implementation into the Einstein Toolkit

Implementation as Einstein Toolkit thorn:

- The Einstein Toolkit is an [open-source modular code optimised for numerical relativity](#) with a large community of users.
- Follow [LlamaWaveHyperboloidal thorn](#): interior Cartesian coordinates matched to outer spherical shell.
- Wave equation suitable for spherical [coordinates](#) everywhere, but infrastructure better suited to Cartesian ones.

Future tests in 3D (NRPy+/Einstein Toolkit):

- [Schwarzschild and Kerr trumpet backgrounds](#) for the scalar field.
- Implement [Good-Bad-Ugly \(GBU\) model](#) ([Gasperin et al. \*Class. Quant. Grav.\* 37.3 \(2020\)](#)), semilinear model of the Einstein eqs.
- [Linerised Einstein equations](#) before non-linear implementation.



# Summary

- Successful **spherically symmetric results** of the hyperboloidal initial value problem → start long-awaited **3D implementation**.
- Current 3D implementation of the **massless wave equation** on a regular hyperboloidal background used as **toy model for GR**.
- **More numerical experiments** (other hyperboloidal backgrounds and models) before attempting the full Einstein equations.

Thank you for your attention!

*Questions?*