Covariant formulation of Entropic Forces: Cosmic Acceleration from First Principles

with Llorenç Espinosa, arXiv:2106.16012 and 16014

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Forces in Physics

• Fundamental Forces

Gravitation, Strong, Weak, E.M.

Residual Forces

Molecular, Nuclear, Surface Tension

Collective Forces

Brownian motion, Entropic Forces

Our proposal

- Entropic Forces are responsible for present cosmic acceleration and many other LSS phenomena.
- Need a covariant formalism of out-of-equilibrium phenomena in GR.
- Just Quantum Mechanics (QFT), (Non eq.) Thermodynamics and GR.

Motivation

- General Relativity is a time-reversible $(t \rightarrow -t)$ theory.
- Most of the history of the universe is adiabatic (actually in local th. eq.), although there is an arrow of time.
- Irreversible phenomena are not included in GR in a systematic way.

Cosmology

Homogeneous and isotropic universe

$$ds^{2} = -dt^{2} + a(t)^{2} \left(\frac{dr^{2}}{1 - kr^{2}} + r^{2}d\Omega_{2}^{2}\right)$$

Filled with a perfect (ideal) fluid

$$T_{\mu\nu} = p g_{\mu\nu} + (\rho + p) u_{\mu} u_{\nu}$$

Friedmann (Einstein) equations

$$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3}\rho \quad (F1) \qquad \qquad \frac{\ddot{a}}{a} = -\frac{4\pi G}{3}\left(\rho + 3p\right) \quad (F2)$$

Cosmology

Covariant Energy-Momentum Tensor Conservation

$$D_{\mu}T^{\mu\nu} = 0 \quad \Rightarrow \quad \dot{\rho} + 3\frac{\dot{a}}{a}(\rho+p) = 0$$

Also derived from Second Law Thermodynamics (in eq.)

$$T\frac{dS}{dt} = \frac{d}{dt}\left(\rho a^{3}\right) + p\frac{d}{dt}\left(a^{3}\right) = 0$$

On a few occasions (e.g. Big Bang)

$$T\frac{dS}{dt} \ge 0$$
 entropy production

Cosmology

Beyond adiabatic cosmology

$$\dot{\rho} + 3\frac{\dot{a}}{a}(\rho + p) = \frac{T\dot{S}}{a^3}$$

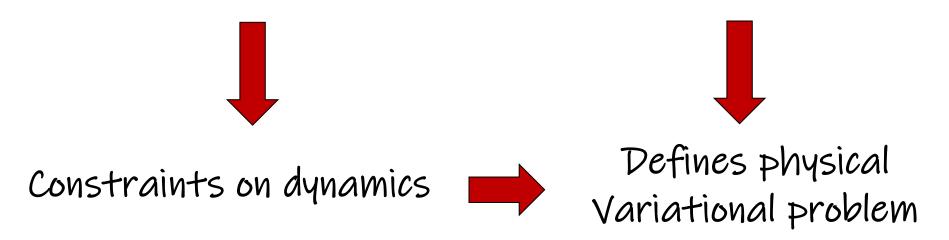
Together with Friedmann equation (F1)

$$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3}\rho$$
 entropic force
$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}\left(\rho + 3p - \frac{T\dot{S}}{a^2\dot{a}}\right)$$

Entropic forces in mechanics

Variational formulation of non-equilibrium thermodynamics (Gay-Balmaz & Yoshimura, 2017)

Laws Thermodynamics + Stationary action principle



Entropic forces in mechanics

General mechanical system with canonical coord. q

$$\mathcal{S} = \int dt \, L\left(q, \, \dot{q}, \, S\right) \quad \blacklozenge \quad \delta \mathcal{S} = \int dt \, \left(\left(\frac{\partial L}{\partial q} - \frac{d}{dt} \frac{\partial L}{\partial \dot{q}} \right) \delta q + \frac{\partial L}{\partial S} \, \delta S \right)$$

+ thermodynamic constraint:

$$\frac{\partial L}{\partial S} \,\delta S = f(q, \, \dot{q}) \,\delta q$$

Constrained equations of motion

 $\frac{d}{dt}\frac{\partial L}{\partial \dot{q}} - \frac{\partial L}{\partial q} = f(q, \dot{q}) \quad \longleftarrow \quad \text{Friction or entropic force}$

Entropic forces in mechanics

what is the physics of $f(q, \dot{q})$?

Has to be obtained from the thermodynamic constraint

$$\frac{\partial L}{\partial S} \dot{S} = f(q, \dot{q}) \dot{q}$$
Usually $\frac{\partial L}{\partial S} = -T \Rightarrow f(q, \dot{q}) \dot{q} \leq 0$ entropy production
Dynamics + Symmetry under
is modified + time-reversal is broken

$$\mathcal{S} = \frac{1}{2\kappa} \int d^4x \sqrt{-g} R + \int d^4x \mathcal{L}_m(g_{\mu\nu}, S)$$

straightforward generalization

$$\delta \mathcal{S} = \int d^4x \left(\frac{1}{2\kappa} \frac{\delta(\sqrt{-gR})}{\delta g^{\mu\nu}} + \frac{\delta \mathcal{L}_m}{\delta g^{\mu\nu}} \right) \delta g^{\mu\nu} + \int d^4x \frac{\partial \mathcal{L}_m}{\partial S} \delta S = 0$$

Variational constraint: 2nd law thermodynamics

$$\frac{\partial \mathcal{L}_m}{\partial s} \delta s = \frac{1}{2} \sqrt{-g} f_{\mu\nu} \delta g^{\mu\nu}$$

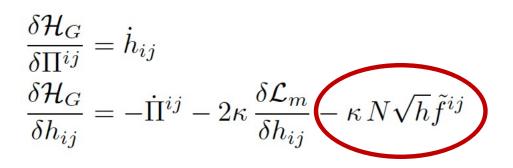
Non-equilibrium Einstein field equations

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = \kappa \left(T_{\mu\nu} - f_{\mu\nu}\right)$$

$$D^{\mu}T_{\mu\nu} = D^{\mu}f_{\mu\nu}$$
entropic force
Bianchi identities

$$\begin{array}{ll} \textbf{ADM Formalism} \\ (3+1) \text{-splitting of space-time} \\ ds^2 &= -(Ndt)^2 + h_{ij}(dx^i + N^i dt)(dx^j + N^j dt) \\ \text{Lie derivative } \pounds_n \text{ along the normal vector } n \\ \text{as a generalization of the time derivative} \\ K_{ij} &= \frac{1}{2}\pounds_n h_{ij} = \frac{1}{2N} \left(\partial_0 h_{ij} - \nabla_i N_j - \nabla_j N_i \right) \\ &= \frac{\partial \mathcal{L}}{\partial s}\pounds_n s = \frac{1}{2}N\sqrt{h} \tilde{f}_{ij}\pounds_n h^{ij} \qquad \tilde{f}_{ij} = h_i^{\mu} h_j^{\nu} f_{\mu\nu} \end{array}$$

Hamilton equations Hamiltonian and momentum constraints



$$\frac{\delta \mathcal{H}_G}{\delta N} = \mathcal{H} = 2\kappa \frac{\partial \mathcal{L}_m}{\partial N}$$
$$\frac{\delta \mathcal{H}_G}{\delta N_i} = \mathcal{H}^i = 2\kappa \frac{\partial \mathcal{L}_m}{\partial N_i}$$

Raychaudhuri Eq.

Congruence of worldlines

$$\Theta_{\mu\nu} = D_{\nu}n_{\mu} = \frac{1}{3}\Theta h_{\mu\nu} + \sigma_{\mu\nu} + \omega_{\mu\nu} - a_{\mu}n_{\nu}$$
$$\pounds_{n}\Theta = -\frac{1}{3}\Theta^{2} - \sigma_{\mu\nu}\sigma^{\mu\nu} + \omega_{\mu\nu}\omega^{\mu\nu} - R_{\mu\nu}n^{\mu}n^{\nu} + D_{\mu}a^{\mu}$$
$$\sigma_{\mu\nu}\sigma^{\mu\nu} > 0 \text{ and } \Theta^{2} > 0 \quad \text{vorticity vanishes } \omega_{\mu\nu} = 0$$
$$R_{\mu\nu}n^{\mu}n^{\nu} = 8\pi G \left(T_{\mu\nu}n^{\mu}n^{\nu} + \frac{1}{2}T - f_{\mu\nu}n^{\mu}n^{\nu} - \frac{1}{2}f\right)$$

If the strong energy condition is satisfied, then: $T_{\mu\nu}n^{\mu}n^{\nu} \ge -\frac{1}{2}T$ and, in the absence of intrinsic acceleration, $a_{\mu} = 0$, we can establish the bound:

$$\pounds_n \Theta + \frac{1}{3} \Theta^2 \le 8\pi G \left(f_{\mu\nu} n^\mu n^\nu + \frac{1}{2} f \right)$$

A positive & suff. large entropic contribution can avoid recollapse

Temperature and Entropy from the matter content

• Mechanical system

 $L(q, \dot{q}, S) = K(q, \dot{q}) - U(q, S) \quad \Rightarrow \quad T = -\frac{\partial L}{\partial S} = \frac{\partial U}{\partial S}$

• Hydrodynamical matter

$$\mathcal{L} = -\sqrt{-g}\,\rho(g_{\mu\nu},\,s) \quad \Rightarrow \quad T = -\frac{1}{\sqrt{-g}}\frac{\partial\mathcal{L}}{\partial s} = -\frac{\partial\rho}{\partial s}$$

Temperature and Entropy from the gravity sector

• Horizon H with induced metric h

$$\mathcal{S}_{\rm GHY} = \frac{1}{8\pi G} \int_H d^3 y \sqrt{h} \, K = \frac{1}{8\pi G} \int_H dt \, \sin\theta d\theta \, d\phi \, \sqrt{h} \, K$$

• Schwarzschild black hole

$$ds^{2} = -\left(1 - \frac{2GM}{r}\right)dt^{2} + \left(1 - \frac{2GM}{r}\right)^{-1}dr^{2} + r^{2}d\Omega_{2}^{2}$$

$$n = -\sqrt{1 - \frac{2GM}{r}}\partial_{r}$$

$$normal \ \text{vector tor}$$

$$S_{2} \ \text{of radius r}$$

$$\begin{split} \mathcal{S}_{\rm GHY} &= \frac{1}{8\pi G} \int_{H} d^{3}y \sqrt{h} \, K = \frac{1}{8\pi G} \int_{H} dt \, \sin\theta d\theta \, d\phi \, \sqrt{h} \, K \\ \sqrt{h}K &= (3GM - 2r) \sin\theta \quad \text{ as event horizon } r = 2GM \\ \mathcal{S}_{\rm GHY} &= -\frac{1}{2} \int dt \, Mc^{2} = -\int dt \, T_{\rm BH} S_{\rm BH} \\ T_{\rm BH} &= \frac{\hbar c^{3}}{8\pi GM} \quad \text{Classical (emergent)} \\ quantum \text{ origin} \\ S_{\rm BH} &= \frac{A \, c^{3}}{4G\hbar} = \frac{4\pi GM^{2}}{\hbar c} \end{split}$$

• Contribution to bulk entropy of the inevitable Schwarzschild black hole component of Dark Matter

Assuming BH total comoving number is conserved, their Total energy density and entropy density $(\hbar = c = 1)$ is

$$\rho_{BH} = n_{BH} M, \quad s_{BH} = n_{BH} 4\pi G M^2$$

Therefore $a^3 \frac{d}{dt} (\rho_{BH} a^3) = T_{BH} \frac{d}{dt} (s_{BH} a^3) = 0$

No contribution to entropic force of the universe unless multiple **black hole mergers** or significant **mass accretion**, which may change **mass or number density** of black holes.

Non-equilibrium thermodynamics in expanding universe

$$ds^{2} = -N(t)^{2}dt^{2} + a^{2}(t)\left(\frac{dr^{2}}{1 - kr^{2}} + r^{2}d\Omega_{2}^{2}\right)$$

$$T^{\mu\nu} = (\rho + p)u^{\mu}u^{\nu} + pg^{\mu\nu} \qquad D^{\mu}T_{\mu\nu} = D^{\mu}f_{\mu\nu}$$
2nd law thermodynamics
$$TdS = d(\rho a^{3}) + p d(a^{3}) \qquad \dot{\rho} + 3H(\rho + p) = \frac{T\dot{S}}{a^{3}}$$

Hamiltonian constraint $\dot{a}^2 + k = \frac{3\pi G}{3}\rho a^2$

Friedmann/Raychaudhuri equation

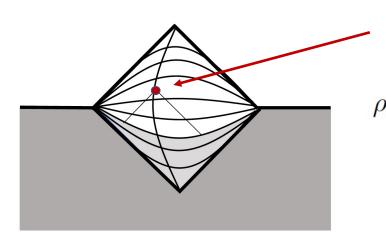
$$\frac{\ddot{a}}{a}=-\frac{4\pi G}{3}(\rho+3p)+\frac{4\pi G}{3}\frac{T\dot{S}}{a^{3}H}$$

• Apparent Cosmological Horizon H

$$n = g^{rr} \partial_r = a^{-1} \sqrt{1 - kr^2} \partial_r$$
unit normal vector to
comoving sphere of radius r
 $\sqrt{h}K = 2N(t) r a \sqrt{1 - kr^2} \sin \theta$ Trace extrinsic curvature
 $r_{AH} = 1/\sqrt{H^2 - k^2/a^2}$ Apparent horizon distance
 $S_{GHY} = -\frac{1}{2G} \int dt N(t) H r_{AH}^2 = -\int dt N(t) T_{AH} S_{AH}$
 $T_{AH} = \frac{\hbar c H}{2\pi}, \quad S_{AH} = \frac{c^3}{\hbar} \frac{\pi r_{AH}^2}{G}$ Emergent

Not enough to contribute to accelerated expansion universe

• Causal Cosmological Horizon H $\sqrt{h}K = 2N(t) r a \sqrt{1 - kr^2 \sin \theta}$ Trace extrinsic curvature $d_H = a \eta$ Causal horizon distance $r = \sinh(\eta \sqrt{-k})/\sqrt{-k}$ Conformal time η $S_{GHY} = -\frac{1}{2G} \int dt N(t) \frac{a}{\sqrt{-k}} \sinh(2\eta\sqrt{-k})$ $= -\int dt N(t) T_H S_H = -\int dt N a^3 \rho_H$ $T_H = \frac{\hbar c}{2\pi} \frac{\sinh(2\eta\sqrt{-k})}{a\eta^2\sqrt{-k}}, \qquad S_H = \frac{\pi c^3}{\hbar} \frac{a^2\eta^2}{G} \qquad \text{Emergent}$



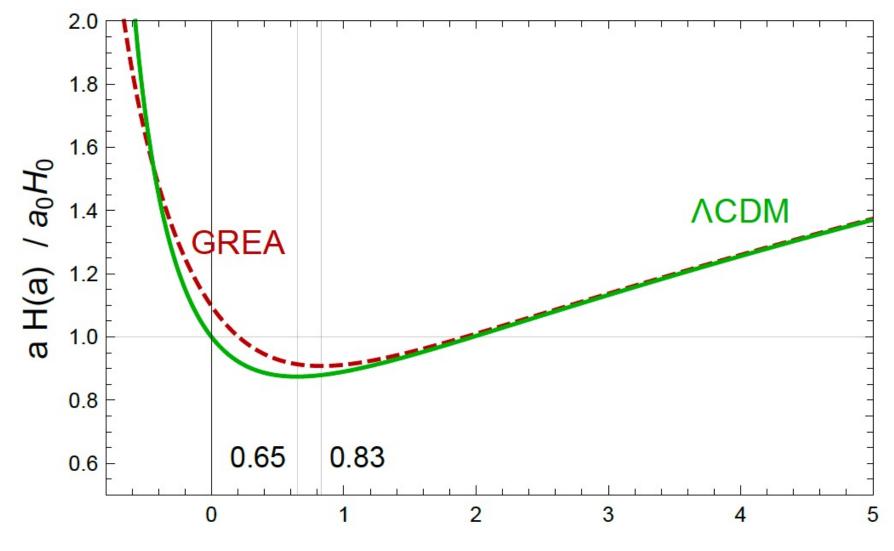
$$v_H a^2 = \frac{T_H S_H}{a} = \frac{x_0}{2G} \sinh(2a_0 H_0 \eta),$$

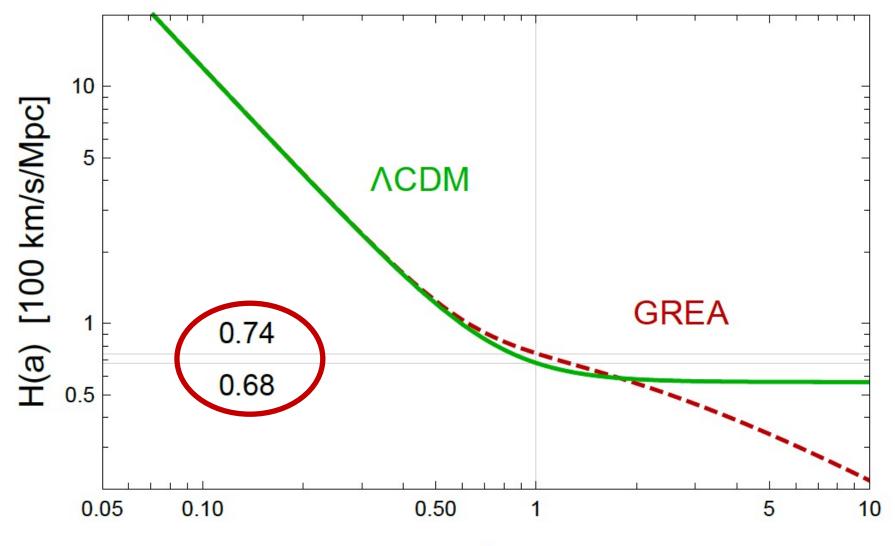
 $x_0 \equiv \frac{1 - \Omega_0}{\Omega_0} = e^{-2N} \left(\frac{T_{\rm rh}}{T_{\rm eq}}\right)^2 (1 + z_{\rm eq}).$

Hamiltonian constraint in conformal time (primes denote derivatives w.r.t. $\tau = a_0 H_0 \eta$)

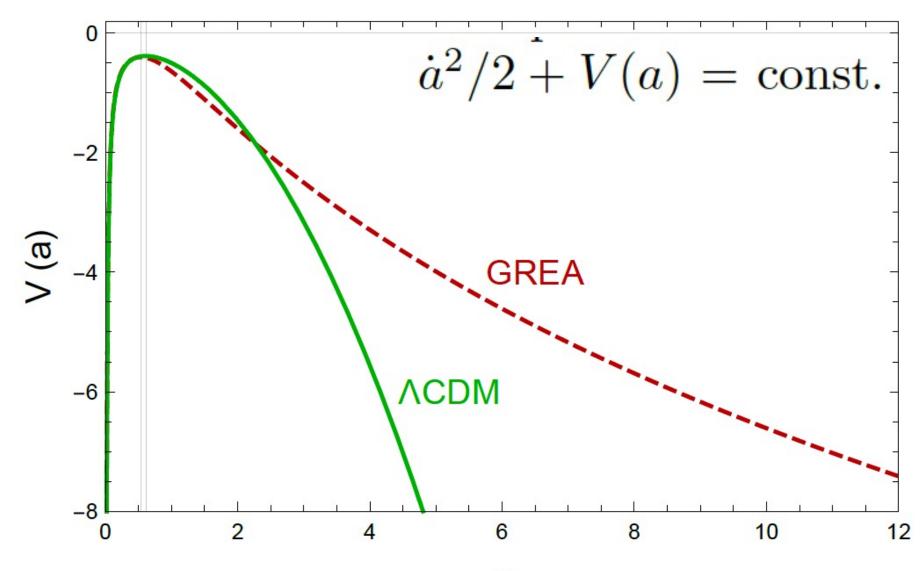
$$\left(\frac{a'}{a_0}\right)^2 = \Omega_{\rm M} \left(\frac{a}{a_0}\right) + \Omega_{\rm K} \left(\frac{a}{a_0}\right)^2 + \left(\frac{4\pi}{3}\Omega_{\rm K} \left(\frac{a}{a_0}\right)^2 \sinh(2\tau)\right)$$

Entropic force term Note: $\Lambda = 0$

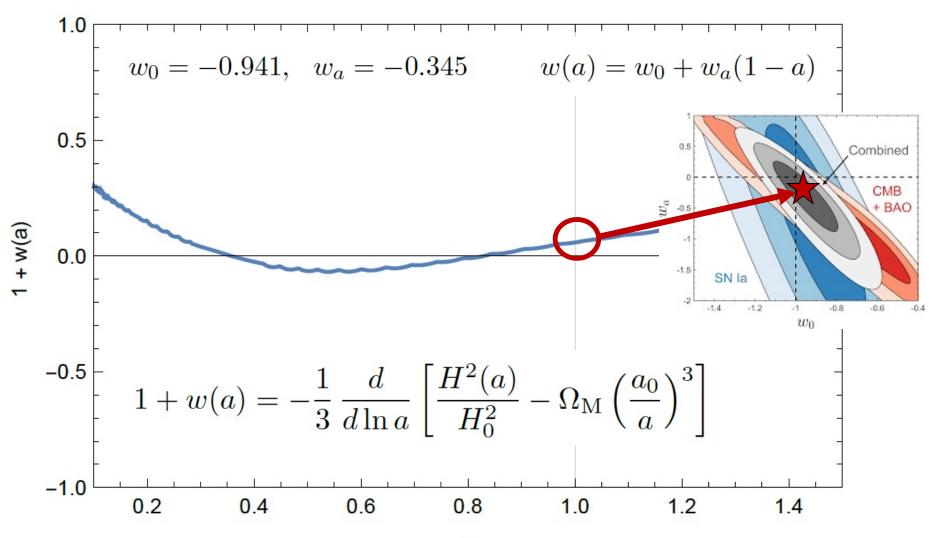




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Conclusions

- Non-equilibrium phenomena in GR: entropic forces
- ADM (3+1) slicing: Raychaudhuri eq. grav. collapse
- Cosmic acceleration from first principles
- No need for a Cosmological Constant
- Just QFT, GR and Non eq. Thermodynamics
- Multiple consequences for Large Scale Structure
- Possible solution of the H_0 , S_8 tensions
- <u>Future</u>: Preheating after inflation (Big Bang)
- <u>Future</u>: Connection w/ Verlinde's emergent gravity
- <u>Future</u>: Connection w/ Buchert's backreaction prob.