

Quantum expectation values on black hole space-times

Elizabeth Winstanley

Consortium for Fundamental Physics
School of Mathematics and Statistics
The University of Sheffield



The
University
Of
Sheffield.

Quantum fields on curved space-time

QFT on curved space-time

QFT on curved space-time

- Classical background space-time
- Quantum field on this background

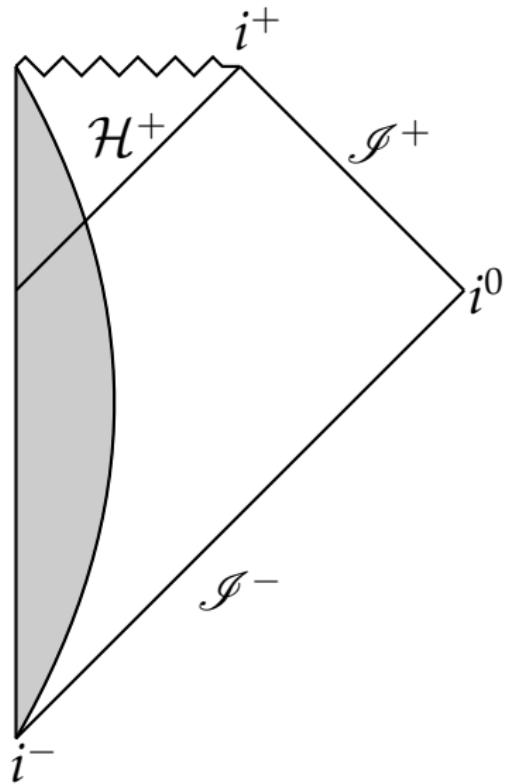
QFT on curved space-time

- Classical background space-time
- Quantum field on this background

Hawking radiation

- Black hole formed by gravitational collapse
- Thermal flux at \mathcal{I}^+

$$T_H = \frac{\kappa}{2\pi}$$



Stress-energy tensor expectation value

Semi-classical Einstein equations

$$G_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi \langle \hat{T}_{\mu\nu} \rangle$$

$c^{-2} \cdot (\text{energy density})$	momentum density			
T^{00}	T^{01}	T^{02}	T^{03}	
T^{10}	T^{11}	T^{12}	T^{13}	
T^{20}	T^{21}	T^{22}	T^{23}	
T^{30}	T^{31}	T^{32}	T^{33}	

Diagram illustrating the components of the stress-energy tensor $T^{\mu\nu}$ in a 4x4 matrix form:

- Rows:**
 - Row 0: T^{00} (red), T^{01} , T^{02} , T^{03}
 - Row 1: T^{10} (orange), T^{11} , T^{12} , T^{13}
 - Row 2: T^{20} (yellow), T^{21} , T^{22} , T^{23}
 - Row 3: T^{30} (blue), T^{31} , T^{32} , T^{33}
- Columns:**
 - Column 0: T^{00} (red), T^{10} (orange), T^{20} (yellow), T^{30} (blue)
 - Column 1: T^{01} , T^{11} , T^{21} , T^{31}
 - Column 2: T^{02} , T^{12} , T^{22} , T^{32}
 - Column 3: T^{03} , T^{13} , T^{23} , T^{33}

Annotations indicate the physical interpretation of the components:

- $c^{-2} \cdot (\text{energy density})$ is associated with the diagonal elements $T^{00}, T^{11}, T^{22}, T^{33}$.
- momentum density is associated with the off-diagonal elements.
- energy flux is associated with the first column ($T^{01}, T^{10}, T^{20}, T^{30}$).
- momentum flux is associated with the second row (T^{01}, T^{02}, T^{03}).
- shear stress is associated with the off-diagonal elements in the momentum density row.
- pressure is associated with the off-diagonal elements in the energy density column.

Massless, conformally coupled scalar field Φ

Klein-Gordon equation

$$\left[\square - \frac{1}{6}R \right] \Phi = 0$$

Massless, conformally coupled scalar field Φ

Klein-Gordon equation

$$\left[\square - \frac{1}{6}R \right] \Phi = 0$$

- \square curved space-time Laplacian $\nabla_\mu \nabla^\mu$

Massless, conformally coupled scalar field Φ

Klein-Gordon equation

$$\left[\square - \frac{1}{6} R \right] \Phi = 0$$

- \square curved space-time Laplacian $\nabla_\mu \nabla^\mu$
- R Ricci scalar curvature

Massless, conformally coupled scalar field Φ

Klein-Gordon equation

$$\left[\square - \frac{1}{6}R \right] \Phi = 0$$

- \square curved space-time Laplacian $\nabla_\mu \nabla^\mu$
- R Ricci scalar curvature

Classical stress-energy tensor

$$\begin{aligned} T_{\mu\nu} = & \frac{2}{3}\Phi_{;\mu}\Phi_{;\nu} - \frac{1}{6}g_{\mu\nu}\Phi^{;\alpha}\Phi_{;\alpha} - \frac{1}{3}\Phi\Phi_{;\mu\nu} \\ & + \frac{1}{3}g_{\mu\nu}\Phi\square\Phi + \frac{1}{6}\left(R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R\right)\Phi^2 \end{aligned}$$

Massless, conformally coupled scalar field Φ

Klein-Gordon equation

$$\left[\square - \frac{1}{6}R \right] \Phi = 0$$

- \square curved space-time Laplacian $\nabla_\mu \nabla^\mu$
- R Ricci scalar curvature

Classical stress-energy tensor

$$\begin{aligned} T_{\mu\nu} = & \frac{2}{3}\Phi_{;\mu}\Phi_{;\nu} - \frac{1}{6}g_{\mu\nu}\Phi^{;\alpha}\Phi_{;\alpha} - \frac{1}{3}\Phi\Phi_{;\mu\nu} \\ & + \frac{1}{3}g_{\mu\nu}\Phi\square\Phi + \frac{1}{6}\left(R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R\right)\Phi^2 \end{aligned}$$

Expectation values of field operators

Expectation values of field operators

Vacuum polarization

$$\langle \hat{\Phi}^2(x) \rangle = \lim_{x' \rightarrow x} [-iG_F(x, x')]$$

Expectation values of field operators

Vacuum polarization

$$\langle \hat{\Phi}^2(x) \rangle = \lim_{x' \rightarrow x} [-iG_F(x, x')]$$

Feynman Green's function $G_F(x, x')$

$$\left[\square - \frac{1}{6}R \right] G_F(x, x') = -(-g)^{-\frac{1}{2}} \delta(x - x')$$

Expectation values of field operators

Vacuum polarization

$$\langle \hat{\Phi}^2(x) \rangle = \lim_{x' \rightarrow x} [-iG_F(x, x')]$$

Feynman Green's function $G_F(x, x')$

$$\left[\square - \frac{1}{6}R \right] G_F(x, x') = -(-g)^{-\frac{1}{2}} \delta(x - x')$$

Stress-energy tensor

$$\langle \hat{T}_{\mu\nu}(x) \rangle = \lim_{x' \rightarrow x} [\mathcal{T}_{\mu\nu}(-iG_F(x, x'))]$$

$\mathcal{T}_{\mu\nu}$ second order differential operator

Feynman Green's function

Static, spherically symmetric black hole

$$ds^2 = -f(r) dt^2 + f(r)^{-1} dr^2 + r^2 d\Omega^2$$

Feynman Green's function

Static, spherically symmetric black hole

$$ds^2 = -f(r) dt^2 + f(r)^{-1} dr^2 + r^2 d\Omega^2$$

Scalar field modes

$$\phi_{\omega\ell m} = \frac{1}{r} e^{-i\omega t} \psi_{\omega\ell}(r) Y_{\ell m}(\theta, \varphi)$$

Feynman Green's function

Static, spherically symmetric black hole

$$ds^2 = -f(r) dt^2 + f(r)^{-1} dr^2 + r^2 d\Omega^2$$

Scalar field modes

$$\phi_{\omega\ell m} = \frac{1}{r} e^{-i\omega t} \psi_{\omega\ell}(r) Y_{\ell m}(\theta, \varphi)$$

Feynman Green's function

Static, spherically symmetric black hole

$$ds^2 = -f(r) dt^2 + f(r)^{-1} dr^2 + r^2 d\Omega^2$$

Scalar field modes

$$\phi_{\omega\ell m} = \frac{1}{r} e^{-i\omega t} \psi_{\omega\ell}(r) Y_{\ell m}(\theta, \varphi)$$

Feynman Green's function

Static, spherically symmetric black hole

$$ds^2 = -f(r) dt^2 + f(r)^{-1} dr^2 + r^2 d\Omega^2$$

Scalar field modes

$$\phi_{\omega\ell m} = \frac{1}{r} e^{-i\omega t} \psi_{\omega\ell}(r) Y_{\ell m}(\theta, \varphi)$$

$$-\frac{d^2 \psi_{\omega\ell}}{dr_*^2} + V_\ell(r_*) \psi_{\omega\ell} = \omega^2 \psi_{\omega\ell} \quad \frac{dr_*}{dr} = \frac{1}{f(r)}$$

Feynman Green's function

Static, spherically symmetric black hole

$$ds^2 = -f(r) dt^2 + f(r)^{-1} dr^2 + r^2 d\Omega^2$$

Scalar field modes

$$\phi_{\omega\ell m} = \frac{1}{r} e^{-i\omega t} \psi_{\omega\ell}(r) Y_{\ell m}(\theta, \varphi)$$

$$-\frac{d^2 \psi_{\omega\ell}}{dr_*^2} + V_\ell(r_*) \psi_{\omega\ell} = \omega^2 \psi_{\omega\ell} \quad \frac{dr_*}{dr} = \frac{1}{f(r)}$$

Feynman Green's function

$$-iG_F(x, x') = \int_{\omega=0}^{\infty} d\omega \sum_{\ell=0}^{\infty} (2\ell+1) \frac{e^{i\omega\Delta t}}{rr'} \psi_{\omega\ell}(r) \psi_{\omega\ell}^*(r') P_\ell(\cos\gamma)$$

Feynman Green's function

Static, spherically symmetric black hole

$$ds^2 = -f(r) dt^2 + f(r)^{-1} dr^2 + r^2 d\Omega^2$$

Scalar field modes

$$\phi_{\omega\ell m} = \frac{1}{r} e^{-i\omega t} \psi_{\omega\ell}(r) Y_{\ell m}(\theta, \varphi)$$

$$-\frac{d^2 \psi_{\omega\ell}}{dr_*^2} + V_\ell(r_*) \psi_{\omega\ell} = \omega^2 \psi_{\omega\ell} \quad \frac{dr_*}{dr} = \frac{1}{f(r)}$$

Feynman Green's function

$$-iG_F(x, x') = \int_{\omega=0}^{\infty} d\omega \sum_{\ell=0}^{\infty} (2\ell+1) \frac{e^{i\omega\Delta t}}{rr'} \psi_{\omega\ell}(r) \psi_{\omega\ell}^*(r') P_\ell(\cos\gamma)$$

Renormalization

DeWitt *Phys. Rept.* **19** 295 (1975)

Christensen *PRD* **14** 2490 (1976)

Wald *CMP* **54** 1 (1977)

Christensen *PRD* **17** 946 (1978)

Decanini & Folacci *PRD* **78** 044025 (2008)

Overall strategy

Stress-energy tensor operator $\hat{T}_{\mu\nu}$

- Involves products of field operators at the same space-time point
- Expectation values are divergent

Overall strategy

Stress-energy tensor operator $\hat{T}_{\mu\nu}$

- Involves products of field operators at the same space-time point
- Expectation values are divergent

Regularization by point-splitting

- $\mathcal{T}_{\mu\nu}(-iG_F(x, x'))$ finite for $x' \neq x$
- Divergences as $x' \rightarrow x$ are purely geometric and independent of the quantum state

Overall strategy

Stress-energy tensor operator $\hat{T}_{\mu\nu}$

- Involves products of field operators at the same space-time point
- Expectation values are divergent

Regularization by point-splitting

- $\mathcal{T}_{\mu\nu}(-iG_F(x, x'))$ finite for $x' \neq x$
- Divergences as $x' \rightarrow x$ are purely geometric and independent of the quantum state

Renormalized expectation value

- Subtract off appropriate divergent terms $G_S(x, x')$

$$\langle \hat{T}_{\mu\nu}(x) \rangle_{\text{ren}} = \lim_{x' \rightarrow x} [\mathcal{T}_{\mu\nu} \{ -i [G_F(x, x') - G_S(x, x')] \}]$$

Hadamard renormalization

[Decanini & Folacci *PRD* **78** 044025 (2008)]

Hadamard renormalization

Hadamard parametrix

$$-\mathrm{i}G_{\mathrm{S}}(x, x') = \frac{U(x, x')}{\sigma(x, x')} + V(x, x') \log\left(\frac{\sigma(x, x')}{\ell^2}\right)$$

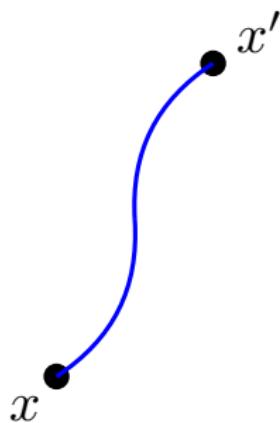
[Decanini & Folacci *PRD* **78** 044025 (2008)]

Hadamard renormalization

Hadamard parametrix

$$-\mathrm{i}G_{\mathrm{S}}(x, x') = \frac{U(x, x')}{\sigma(x, x')} + V(x, x') \log\left(\frac{\sigma(x, x')}{\ell^2}\right)$$

- $2\sigma(x, x')$ square of the geodesic distance between x and x'



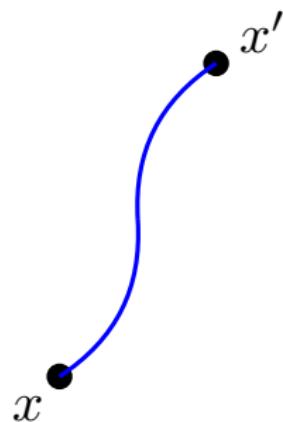
[Decanini & Folacci PRD 78 044025 (2008)]

Hadamard renormalization

Hadamard parametrix

$$-\mathrm{i}G_{\mathrm{S}}(x, x') = \frac{U(x, x')}{\sigma(x, x')} + V(x, x') \log\left(\frac{\sigma(x, x')}{\ell^2}\right)$$

- $2\sigma(x, x')$ square of the geodesic distance between x and x'
- $U(x, x'), V(x, x')$ symmetric biscalars regular as $x' \rightarrow x$



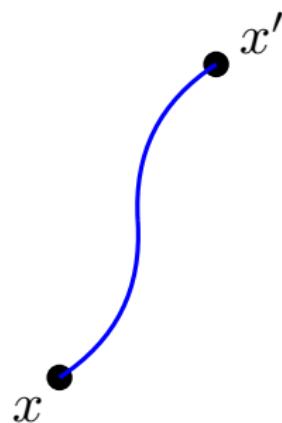
[Decanini & Folacci PRD 78 044025 (2008)]

Hadamard renormalization

Hadamard parametrix

$$-\mathrm{i}G_{\mathrm{S}}(x, x') = \frac{U(x, x')}{\sigma(x, x')} + V(x, x') \log\left(\frac{\sigma(x, x')}{\ell^2}\right)$$

- $2\sigma(x, x')$ square of the geodesic distance between x and x'
- $U(x, x'), V(x, x')$ symmetric biscalars regular as $x' \rightarrow x$
- ℓ renormalization length scale



[Decanini & Folacci PRD 78 044025 (2008)]

Renormalized expectation values

$$\langle \hat{T}_{\mu\nu}(x) \rangle_{\text{ren}} = \lim_{x' \rightarrow x} [\mathcal{T}_{\mu\nu} \{ -i [G_F(x, x') - G_S(x, x')] \}]$$

Renormalized expectation values

$$\langle \hat{T}_{\mu\nu}(x) \rangle_{\text{ren}} = \lim_{x' \rightarrow x} [\mathcal{T}_{\mu\nu} \{ -i [G_F(x, x') - G_S(x, x')] \}]$$

Feynman Green's function $G_F(x, x')$

- Mode sum over separable solutions of the Klein-Gordon equation
- Modes typically found numerically

Renormalized expectation values

$$\langle \hat{T}_{\mu\nu}(x) \rangle_{\text{ren}} = \lim_{x' \rightarrow x} [\mathcal{T}_{\mu\nu} \{ -i [G_F(x, x') - G_S(x, x')] \}]$$

Feynman Green's function $G_F(x, x')$

- Mode sum over separable solutions of the Klein-Gordon equation
- Modes typically found numerically

Hadamard parametrix $G_S(x, x')$

- Purely geometric
- Taylor series expansions for x' close to x

Renormalized expectation values

$$\langle \hat{T}_{\mu\nu}(x) \rangle_{\text{ren}} = \lim_{x' \rightarrow x} [\mathcal{T}_{\mu\nu} \{ -i [G_F(x, x') - G_S(x, x')] \}]$$

Feynman Green's function $G_F(x, x')$

- Mode sum over separable solutions of the Klein-Gordon equation
- Modes typically found numerically

Hadamard parametrix $G_S(x, x')$

- Purely geometric
- Taylor series expansions for x' close to x

Challenge

How can we subtract $G_S(x, x')$ from $G_F(x, x')$ so that limit can be taken and answer computed numerically?

WKB-based method

Candelas & Howard *PRD* **29** 1618 (1984)

Howard & Candelas *PRL* **53** 403 (1984)

Howard *PRD* **30** 2532 (1984)

Anderson, Hiscock & Samuel *PRD* **51** 4337 (1995)

EW & Young *PRD* **77** 024008 (2008)

Flachi & Tanaka *PRD* **78** 064011 (2008)

Breen & Ottewill *PRD* **82** 084019 (2010)

Breen & Ottewill *PRD* **85** 084029 (2012)

WKB-based method

$$\langle \hat{\Phi}^2(x) \rangle_{\text{ren}} = \lim_{x' \rightarrow x} \left\{ -i [G_F(x, x') - G_S(x, x')] \right\}$$

WKB-based method

$$\langle \hat{\Phi}^2(x) \rangle_{\text{ren}} = \lim_{x' \rightarrow x} \left\{ -i [G_F(x, x') - G_S(x, x')] \right\}$$

Time-like point splitting

$$-iG_F(x, x') = \int_{\omega=0}^{\infty} d\omega \sum_{\ell=0}^{\infty} (2\ell + 1) \frac{e^{i\omega\epsilon}}{r^2} |\psi_{\omega\ell}(r)|^2$$

WKB-based method

$$\langle \hat{\Phi}^2(x) \rangle_{\text{ren}} = \lim_{x' \rightarrow x} \left\{ -i [G_F(x, x') - G_S(x, x')] \right\}$$

Time-like point splitting

$$-iG_F(x, x') = \int_{\omega=0}^{\infty} d\omega \sum_{\ell=0}^{\infty} (2\ell + 1) \frac{e^{i\omega\epsilon}}{r^2} |\psi_{\omega\ell}(r)|^2$$

$$\begin{aligned} \langle \hat{\Phi}^2(x) \rangle_{\text{ren}} = & \lim_{\epsilon \rightarrow 0} \left\{ \int_{\omega=0}^{\infty} d\omega \sum_{\ell=0}^{\infty} \left[(2\ell + 1) \frac{e^{i\omega\epsilon}}{r^2} |\psi_{\omega\ell}(r)|^2 \right. \right. \\ & \left. \left. + iG_S(\epsilon) \right] \right\} \end{aligned}$$

WKB-based method

$$\langle \hat{\Phi}^2(x) \rangle_{\text{ren}} = \lim_{x' \rightarrow x} \left\{ -i [G_F(x, x') - G_S(x, x')] \right\}$$

Time-like point splitting

$$-iG_F(x, x') = \int_{\omega=0}^{\infty} d\omega \sum_{\ell=0}^{\infty} (2\ell+1) \frac{e^{i\omega\epsilon}}{r^2} |\psi_{\omega\ell}(r)|^2$$

$$\begin{aligned} \langle \hat{\Phi}^2(x) \rangle_{\text{ren}} &= \lim_{\epsilon \rightarrow 0} \left\{ \int_{\omega=0}^{\infty} d\omega \sum_{\ell=0}^{\infty} \left[(2\ell+1) \frac{e^{i\omega\epsilon}}{r^2} |\psi_{\omega\ell}(r)|^2 - \text{WKB} \right] \right. \\ &\quad \left. + \left[\int_{\omega=0}^{\infty} d\omega \sum_{\ell=0}^{\infty} \text{WKB} \right] + iG_S(\epsilon) \right\} \end{aligned}$$

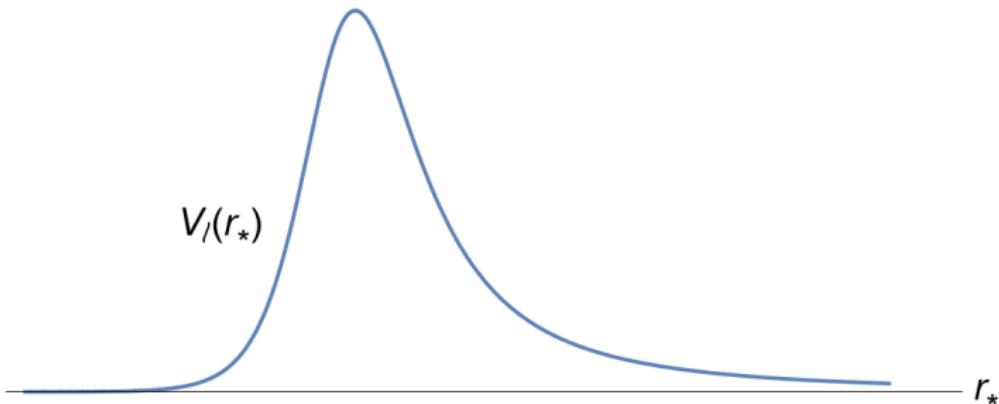
WKB approximation

WKB approximation

$$-\frac{d^2\psi_{\omega\ell}}{dr_*^2} + V_\ell(r_*)\psi_{\omega\ell} = \omega^2\psi_{\omega\ell}$$

WKB approximation

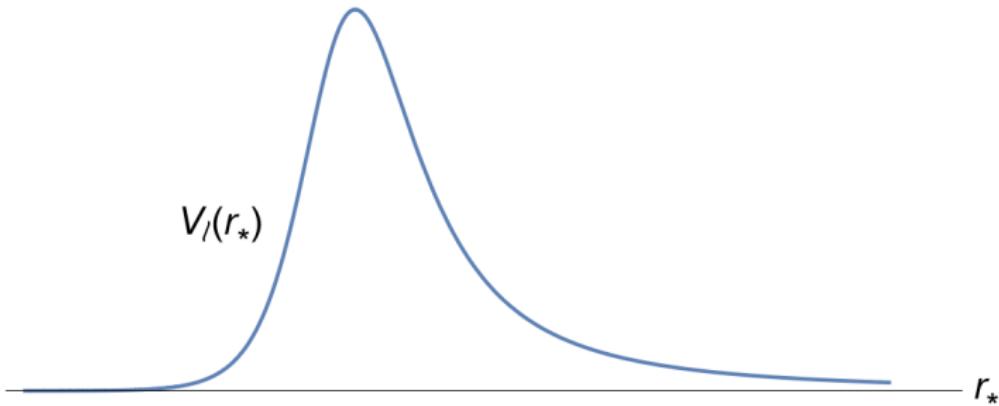
$$-\frac{d^2\psi_{\omega\ell}}{dr_*^2} + V_\ell(r_*)\psi_{\omega\ell} = \omega^2\psi_{\omega\ell}$$



WKB approximation

$$-\frac{d^2\psi_{\omega\ell}}{dr_*^2} + V_\ell(r_*)\psi_{\omega\ell} = \omega^2\psi_{\omega\ell}$$

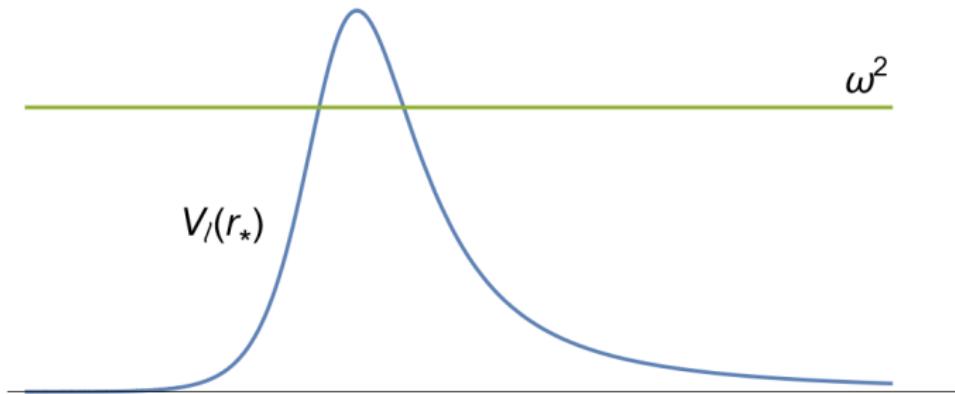
$$r_* \rightarrow \pm\infty \quad V_\ell(r_*) \rightarrow 0 \quad \psi_{\omega\ell} \sim e^{\pm i\omega r_*}$$



WKB approximation

$$-\frac{d^2\psi_{\omega\ell}}{dr_*^2} + V_\ell(r_*)\psi_{\omega\ell} = \omega^2\psi_{\omega\ell}$$

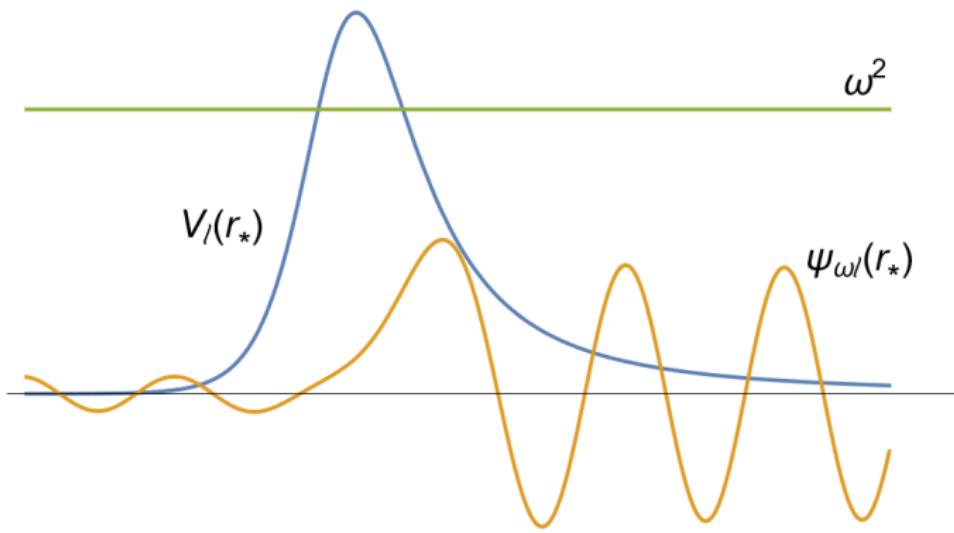
$$r_* \rightarrow \pm\infty \quad V_\ell(r_*) \rightarrow 0 \quad \psi_{\omega\ell} \sim e^{\pm i\omega r_*}$$



WKB approximation

$$-\frac{d^2\psi_{\omega\ell}}{dr_*^2} + V_\ell(r_*)\psi_{\omega\ell} = \omega^2\psi_{\omega\ell}$$

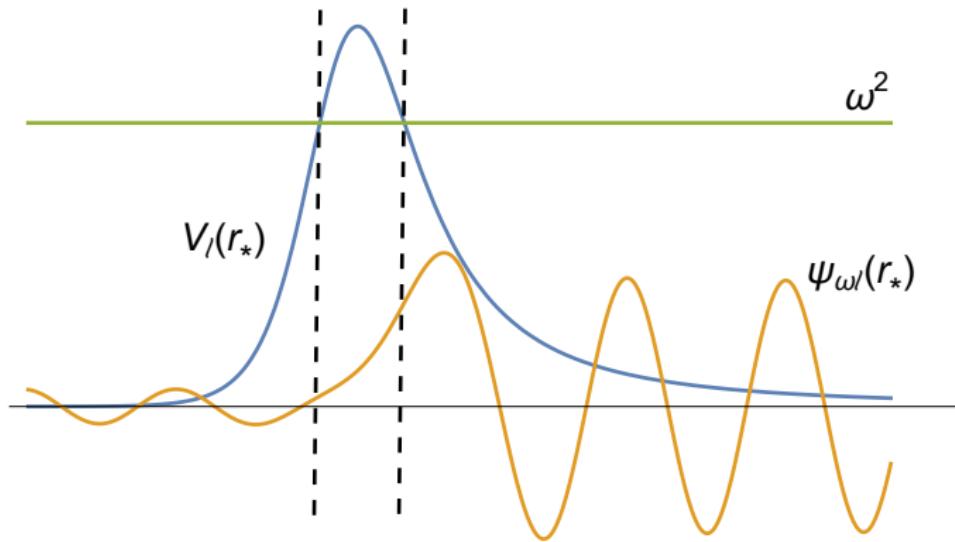
$$r_* \rightarrow \pm\infty \quad V_\ell(r_*) \rightarrow 0 \quad \psi_{\omega\ell} \sim e^{\pm i\omega r_*}$$



WKB approximation

$$-\frac{d^2\psi_{\omega\ell}}{dr_*^2} + V_\ell(r_*)\psi_{\omega\ell} = \omega^2\psi_{\omega\ell}$$

$$r_* \rightarrow \pm\infty \quad V_\ell(r_*) \rightarrow 0 \quad \psi_{\omega\ell} \sim e^{\pm i\omega r_*}$$



Euclideanization

Euclideanization

- Wick rotation $t \rightarrow -i\tau, \omega \rightarrow i\omega$

Euclideanization

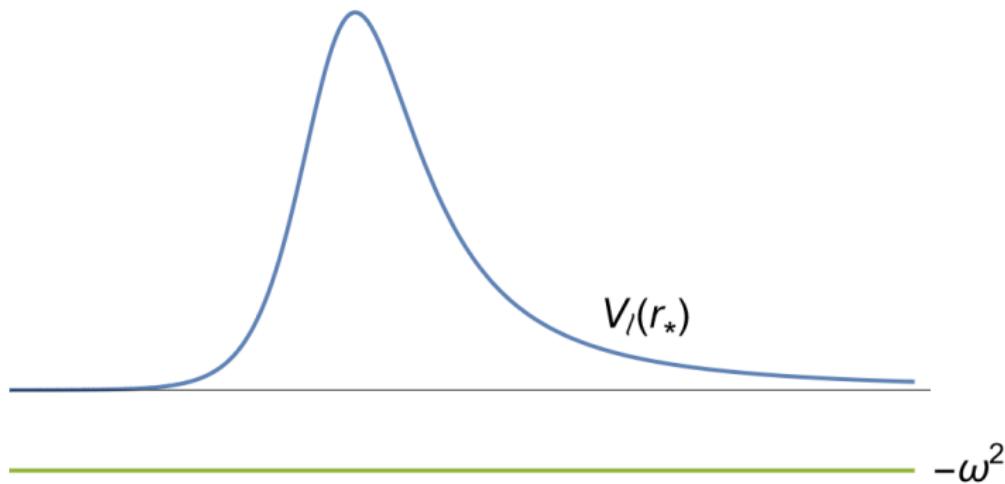
- Wick rotation $t \rightarrow -i\tau, \omega \rightarrow i\omega$
- Radial equation

$$-\frac{d^2\psi_{\omega\ell}}{dr_*^2} + V_\ell(r_*)\psi_{\omega\ell} = -\omega^2\psi_{\omega\ell}$$

Euclideanization

- Wick rotation $t \rightarrow -i\tau, \omega \rightarrow i\omega$
- Radial equation

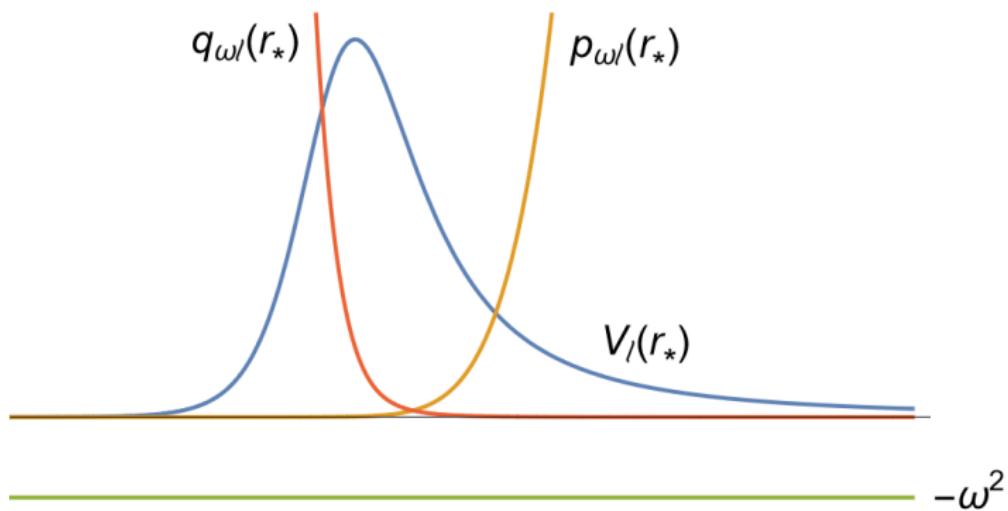
$$-\frac{d^2\psi_{\omega\ell}}{dr_*^2} + V_\ell(r_*)\psi_{\omega\ell} = -\omega^2\psi_{\omega\ell}$$



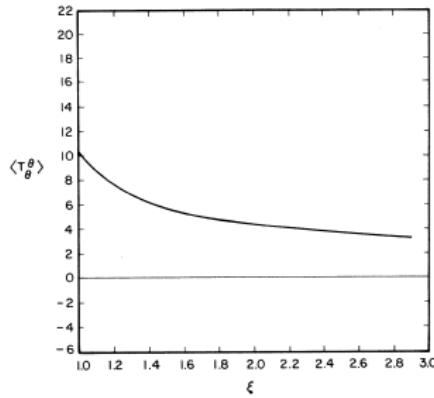
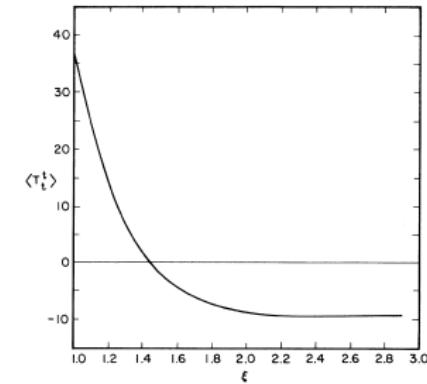
Euclideanization

- Wick rotation $t \rightarrow -i\tau, \omega \rightarrow i\omega$
- Radial equation

$$-\frac{d^2\psi_{\omega\ell}}{dr_*^2} + V_\ell(r_*)\psi_{\omega\ell} = -\omega^2\psi_{\omega\ell}$$

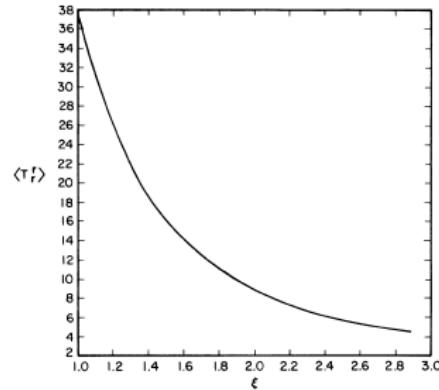


RSET on Schwarzschild



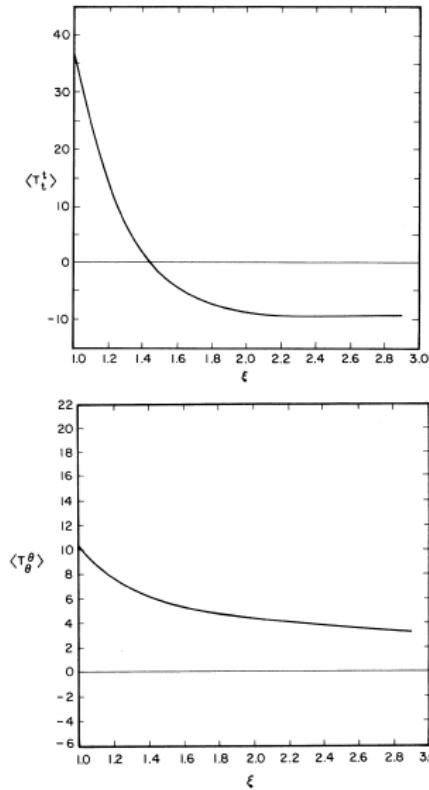
$$ds^2 = f(r) d\tau^2 + f(r)^{-1} dr^2 + r^2 d\Omega^2$$

$$f(r) = 1 - \frac{2M}{r}$$

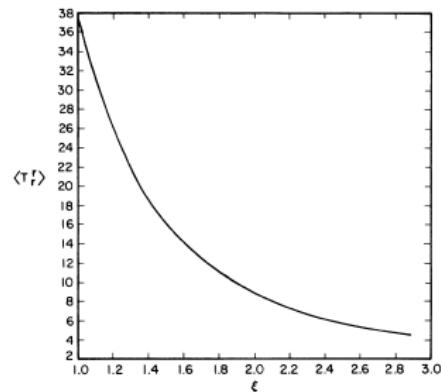


[Howard & Candelas PRL 53 403 (1984)]

RSET on Schwarzschild

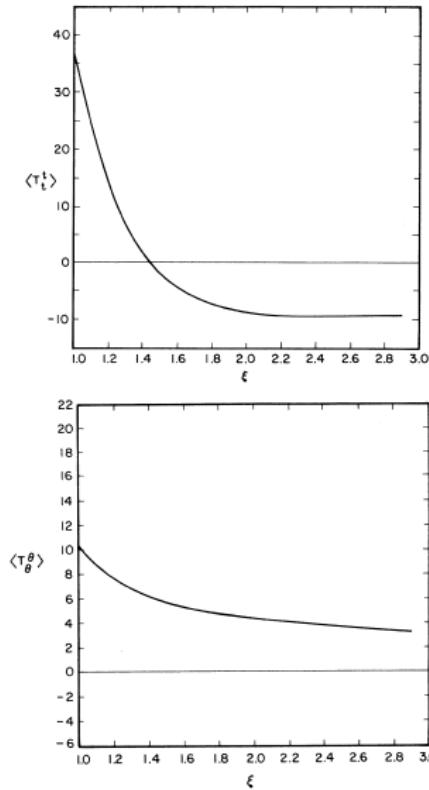


- Euclidean time coordinate periodic $\tau \rightarrow \tau + 2\pi/\kappa$

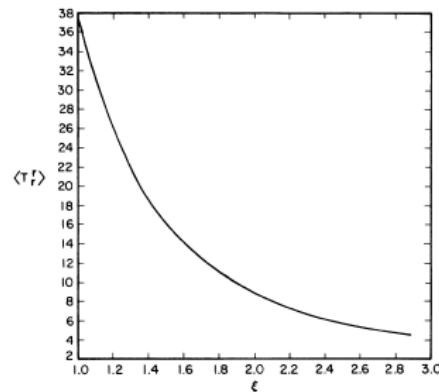


[Howard & Candelas *PRL* **53** 403 (1984)]

RSET on Schwarzschild



- Euclidean time coordinate periodic $\tau \rightarrow \tau + 2\pi/\kappa$
- Thermal state at temperature $\kappa/2\pi$



[Howard & Candelas *PRL* **53** 403 (1984)]

Extended coordinates method

Taylor & Breen *PRD* **94** 125024 (2016)

Taylor & Breen *PRD* **96** 105020 (2017)

Morley, Taylor & EW *CQG* **35** 235010 (2018)

Breen & Taylor *PRD* **98** 105006 (2018)

Morley, Taylor & EW *PRD* **103** 045007 (2021)

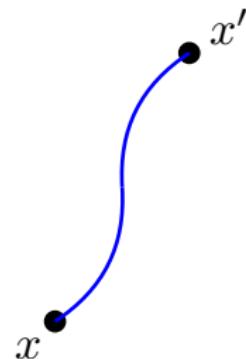
Renormalization

Renormalization

$$\langle \hat{\Phi}^2(x) \rangle_{\text{ren}} = \lim_{x' \rightarrow x} \left\{ -i [G_F(x, x') - G_S(x, x')] \right\}$$

Renormalization

$$\langle \hat{\Phi}^2(x) \rangle_{\text{ren}} = \lim_{x' \rightarrow x} \left\{ -i [G_F(x, x') - G_S(x, x')] \right\}$$

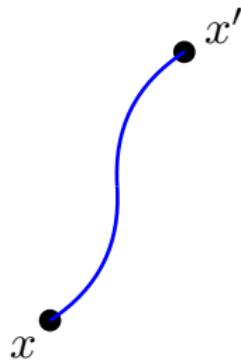
$$-iG_S(x, x') = \frac{U(x, x')}{\sigma(x, x')} + V(x, x') \log \left(\frac{\sigma(x, x')}{\ell^2} \right)$$


Renormalization

$$\langle \hat{\Phi}^2(x) \rangle_{\text{ren}} = \lim_{x' \rightarrow x} \left\{ -i [G_F(x, x') - G_S(x, x')] \right\}$$

$$-iG_S(x, x') = \frac{U(x, x')}{\sigma(x, x')} + V(x, x') \log \left(\frac{\sigma(x, x')}{\ell^2} \right)$$

- $\sigma(x, x') = 0$ for null separated points

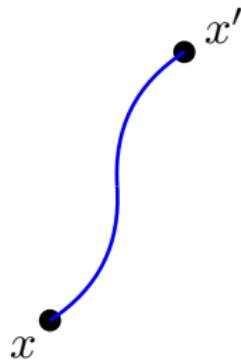


Renormalization

$$\langle \hat{\Phi}^2(x) \rangle_{\text{ren}} = \lim_{x' \rightarrow x} \left\{ -i [G_F(x, x') - G_S(x, x')] \right\}$$

$$-iG_S(x, x') = \frac{U(x, x')}{\sigma(x, x')} + V(x, x') \log \left(\frac{\sigma(x, x')}{\ell^2} \right)$$

- $\sigma(x, x') = 0$ for null separated points
- Euclideanization

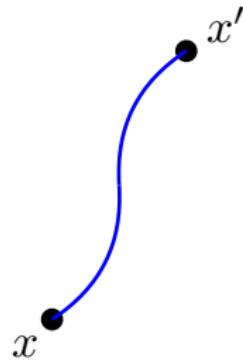


Renormalization

$$\langle \hat{\Phi}^2(x) \rangle_{\text{ren}} = \lim_{x' \rightarrow x} \left\{ -i [G_F(x, x') - G_S(x, x')] \right\}$$

$$-iG_S(x, x') = \frac{U(x, x')}{\sigma(x, x')} + V(x, x') \log \left(\frac{\sigma(x, x')}{\ell^2} \right)$$

- $\sigma(x, x') = 0$ for null separated points
- Euclideanization

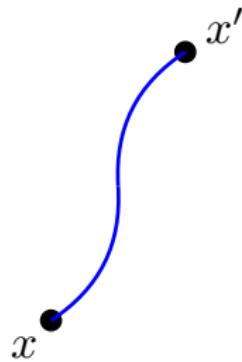


Renormalization

$$\langle \hat{\Phi}^2(x) \rangle_{\text{ren}} = \lim_{x' \rightarrow x} \{ [G_E(x, x') + iG_S(x, x')] \}$$

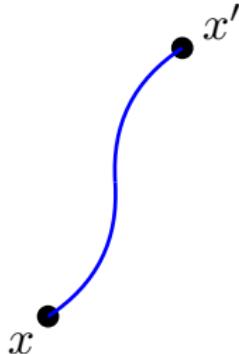
$$-iG_S(x, x') = \frac{U(x, x')}{\sigma(x, x')} + V(x, x') \log \left(\frac{\sigma(x, x')}{\ell^2} \right)$$

- $\sigma(x, x') = 0$ for null separated points
- Euclideanization



Renormalization

$$\langle \hat{\Phi}^2(x) \rangle_{\text{ren}} = \lim_{x' \rightarrow x} \{ [G_E(x, x') + iG_S(x, x')] \}$$

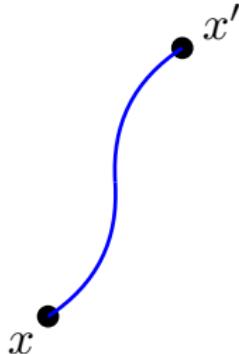
$$-iG_S(x, x') = \frac{U(x, x')}{\sigma(x, x')} + V(x, x') \log \left(\frac{\sigma(x, x')}{\ell^2} \right)$$


- $\sigma(x, x') = 0$ for null separated points
- Euclideanization

$$G_E(x, x') = \frac{\kappa}{8\pi^2} \sum_{n=-\infty}^{\infty} \sum_{\ell=0}^{\infty} e^{ink\Delta\tau} (2\ell+1) p_{n\ell}(r) q_{n\ell}(r) P_\ell(\cos\gamma)$$

Renormalization

$$\langle \hat{\Phi}^2(x) \rangle_{\text{ren}} = \lim_{x' \rightarrow x} \{ [G_E(x, x') + iG_S(x, x')] \}$$

$$-iG_S(x, x') = \frac{U(x, x')}{\sigma(x, x')} + V(x, x') \log \left(\frac{\sigma(x, x')}{\ell^2} \right)$$


- $\sigma(x, x') = 0$ for null separated points
- Euclideanization

$$G_E(x, x') = \frac{\kappa}{8\pi^2} \sum_{n=-\infty}^{\infty} \sum_{\ell=0}^{\infty} e^{ink\Delta\tau} (2\ell + 1) p_{n\ell}(r) q_{n\ell}(r) P_\ell(\cos\gamma)$$

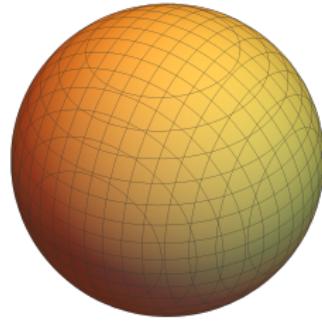
$$-iG_S(x, x') = \frac{\kappa}{8\pi^2} \sum_{n=-\infty}^{\infty} \sum_{\ell=0}^{\infty} e^{ink\Delta\tau} (2\ell + 1) \Gamma_{n\ell}(r) P_\ell(\cos\gamma)$$

Topological black holes

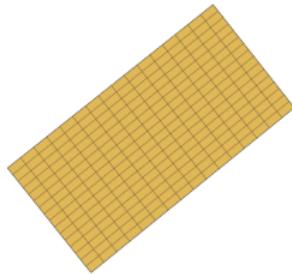
Schwarzschild-anti-de Sitter space-time

$$ds^2 = - \left(k - \frac{2M}{r} + \frac{r^2}{L^2} \right) dt^2 + \left(k - \frac{2M}{r} + \frac{r^2}{L^2} \right)^{-1} dr^2 + r^2 d\Omega_k^2$$

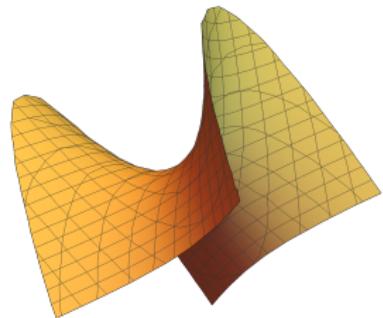
$$k = 1$$



$$k = 0$$

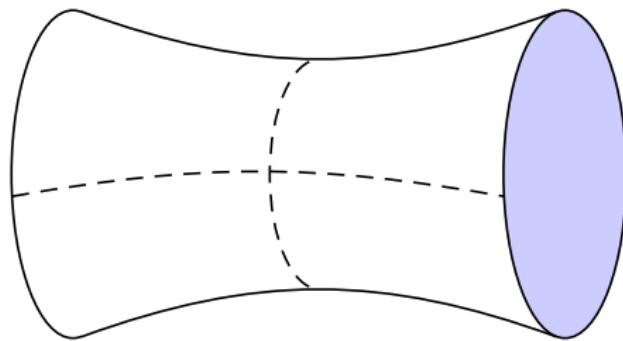


$$k = -1$$



$$d\Omega_1^2 = d\theta^2 + \sin^2 \theta d\varphi^2 \quad d\Omega_0^2 = d\theta^2 + \theta^2 d\varphi^2 \quad d\Omega_{-1}^2 = d\theta^2 + \sinh^2 \theta d\varphi^2$$

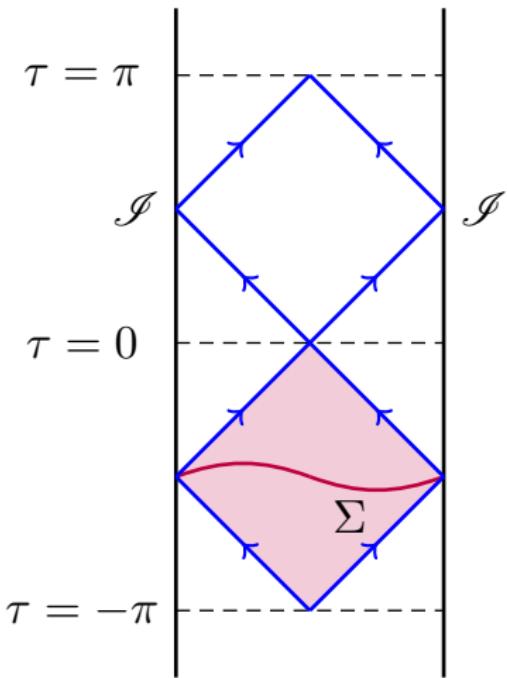
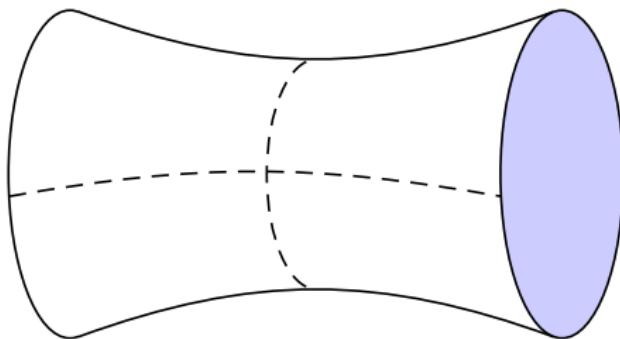
Geometry of anti-de Sitter space-time



- Time-like boundary
- Boundary conditions required

[Avis, Isham & Storey *PRD* **18** 3565 (1978)]

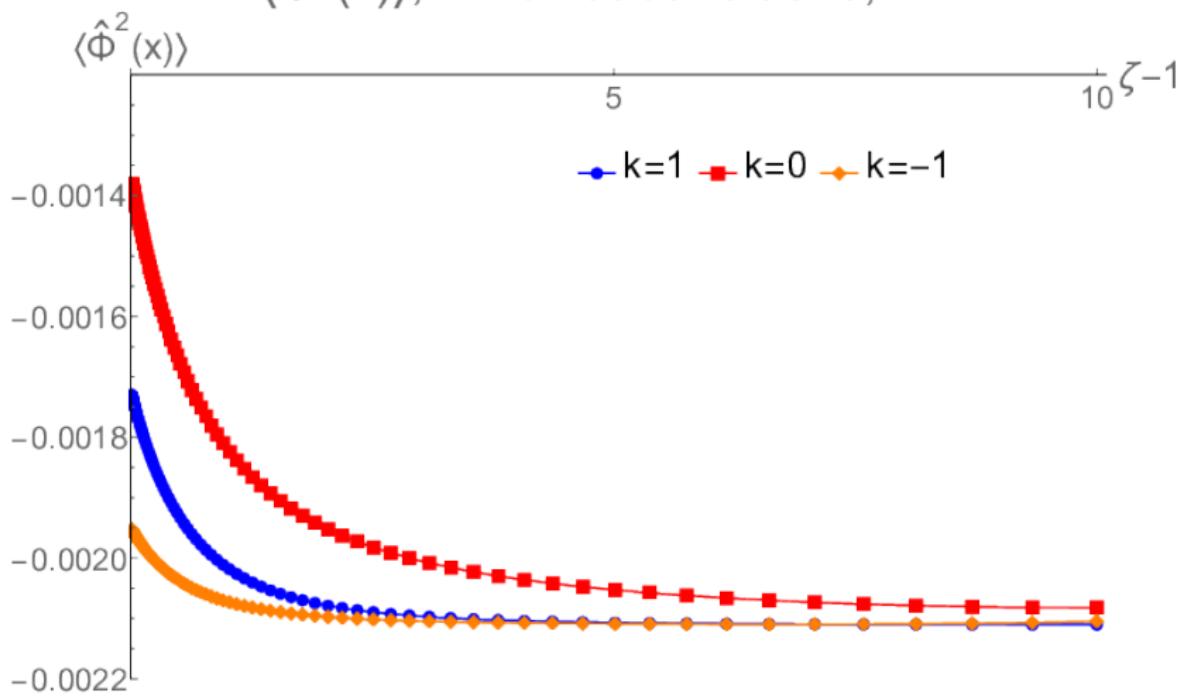
Geometry of anti-de Sitter space-time



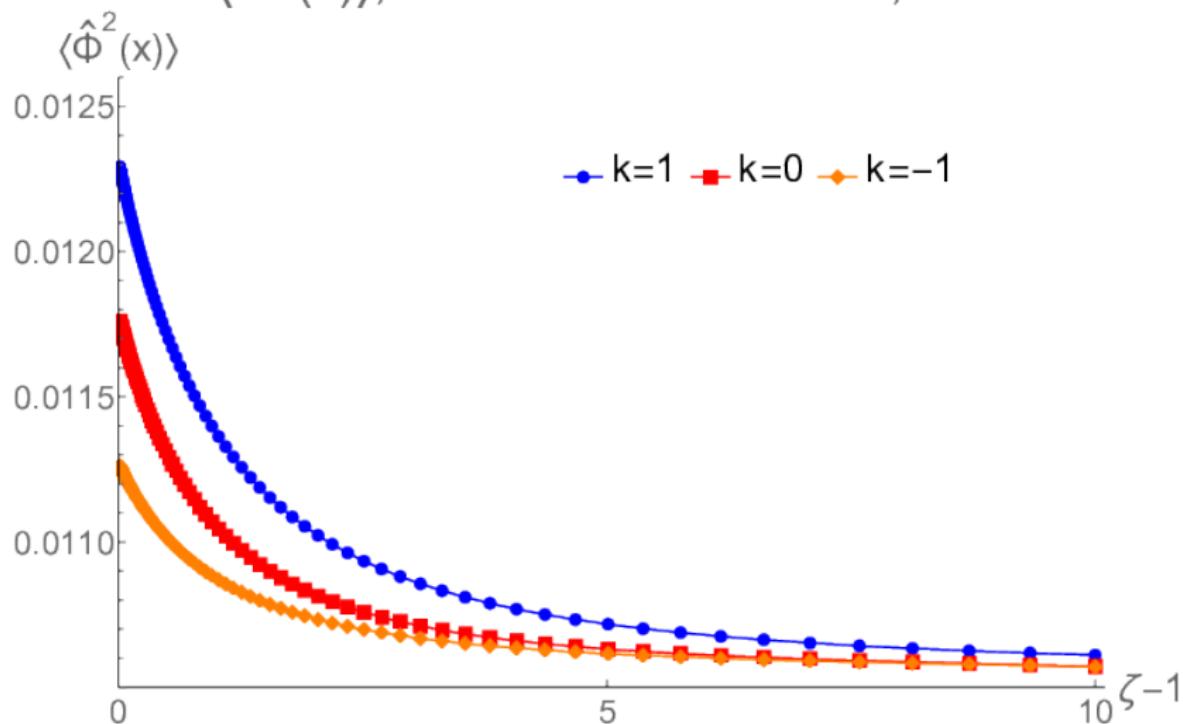
- Time-like boundary
- Boundary conditions required

[Avis, Isham & Storey *PRD* **18** 3565 (1978)]

$\langle \hat{\Phi}^2(x) \rangle$, Dirichlet conditions, $L=1$



[Flachi & Tanaka *PRD* **78** 064011 (2008), Morley, Taylor & EW *CQG* **35** 235010 (2018)]

$\langle \hat{\Phi}^2(x) \rangle$, Neumann conditions, L=1

[Morley, Taylor & EW PRD 103 045007 (2021)]

Pragmatic mode sum method

Levi & Ori *PRD* **91** 104028 (2015)

Levi & Ori *PRD* **94** 044054 (2016)

Levi & Ori *PRL* **117** 231101 (2016)

Levi, Eilon, Ori & van de Meent *PRL* **118** 141102 (2017)

Levi *PRD* **95** 025007 (2017)

Lanir, Levi, Ori & Sela *PRD* **97** 024033 (2018)

Lanir, Levi & Ori *PRD* **98** 084017 (2018)

Lanir, Ori, Zilberman, Sela, Maline & Levi *PRD* **99** 061502 (2019)

Zilberman, Levi & Ori *PRL* **124** 171302 (2020)

Zilberman & Ori *PRD* **104** 024066 (2021)

Lorentzian black hole space-times

Lorentzian black hole space-times

$$\langle \hat{T}_{\mu\nu}(x) \rangle_{\text{ren}} = \lim_{x' \rightarrow x} [\mathcal{T}_{\mu\nu} \{ -i [G_F(x, x') - G_S(x, x')] \}]$$

Lorentzian black hole space-times

$$\langle \hat{T}_{\mu\nu}(x) \rangle_{\text{ren}} = \lim_{x' \rightarrow x} [\mathcal{T}_{\mu\nu} \{ -i [G_F(x, x') - G_S(x, x')] \}]$$

Time-like point splitting

$$-iG_F(x, x') = \int_{\omega=0}^{\infty} d\omega \sum_{\ell=0}^{\infty} (2\ell + 1) \frac{e^{i\omega\epsilon}}{r^2} |\psi_{\omega\ell}(r)|^2$$

Lorentzian black hole space-times

$$\langle \hat{T}_{\mu\nu}(x) \rangle_{\text{ren}} = \lim_{x' \rightarrow x} [\mathcal{T}_{\mu\nu} \{ -i [G_F(x, x') - G_S(x, x')] \}]$$

Time-like point splitting

$$-iG_F(x, x') = \int_{\omega=0}^{\infty} d\omega \sum_{\ell=0}^{\infty} (2\ell + 1) \frac{e^{i\omega\epsilon}}{r^2} |\psi_{\omega\ell}(r)|^2$$

$$-iG_S(x, x') = -\frac{\mathcal{A}(r)}{\epsilon^2} - \mathcal{B}(r) \log\left(\frac{f\epsilon^2}{4\ell^2}\right) + \text{finite terms}$$

Lorentzian black hole space-times

$$\langle \hat{T}_{\mu\nu}(x) \rangle_{\text{ren}} = \lim_{x' \rightarrow x} [\mathcal{T}_{\mu\nu} \{ -i [G_F(x, x') - G_S(x, x')] \}]$$

Time-like point splitting

$$-iG_F(x, x') = \int_{\omega=0}^{\infty} d\omega \sum_{\ell=0}^{\infty} (2\ell + 1) \frac{e^{i\omega\epsilon}}{r^2} |\psi_{\omega\ell}(r)|^2$$

$$\begin{aligned} -iG_S(x, x') &= -\frac{\mathcal{A}(r)}{\epsilon^2} - \mathcal{B}(r) \log\left(\frac{f\epsilon^2}{4\ell^2}\right) + \text{finite terms} \\ &= \int_{\omega=0}^{\infty} \left[\mathcal{A}(r)\omega + \frac{\mathcal{B}(r)}{\omega + \sqrt{f}/2\ell} \right] e^{i\omega\epsilon} d\omega + \mathcal{C}(r) + \dots \end{aligned}$$

Generalized integrals

[Levi & Ori *PRD* **91** 104028 (2015)]

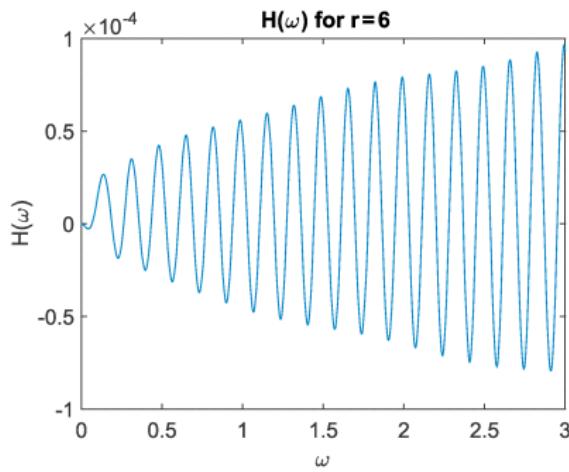
Generalized integrals

$$\mathcal{H}(\omega) = \int_{\omega=0}^{\omega} d\omega e^{i\omega\epsilon} \mathcal{G}_{\omega}(r)$$

[Levi & Ori *PRD* **91** 104028 (2015)]

Generalized integrals

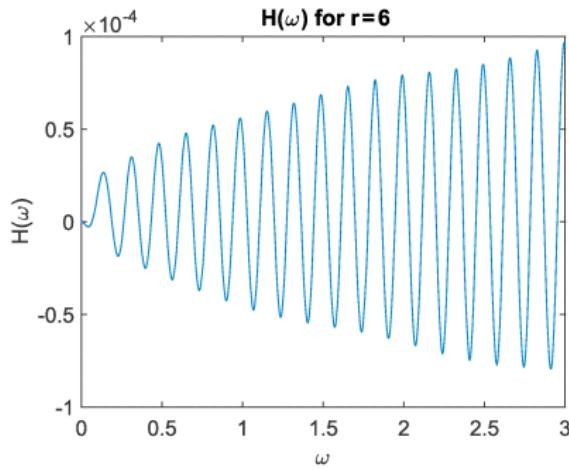
$$\mathcal{H}(\omega) = \int_{\omega=0}^{\omega} d\omega e^{i\omega\epsilon} \mathcal{G}_{\omega}(r)$$



[Levi & Ori *PRD* **91** 104028 (2015)]

Generalized integrals

$$\mathcal{H}(\omega) = \int_{\omega=0}^{\omega} d\omega e^{i\omega\epsilon} \mathcal{G}_{\omega}(r) \quad \mathcal{H}_*(\omega) = \frac{1}{2} [\mathcal{H}(\omega) + \mathcal{H}(\omega + \nu/2)]$$

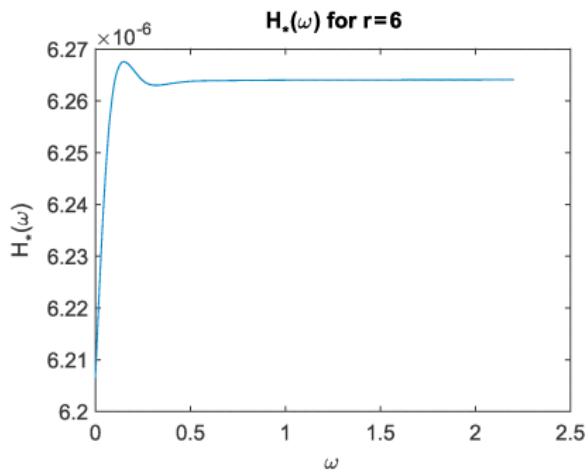
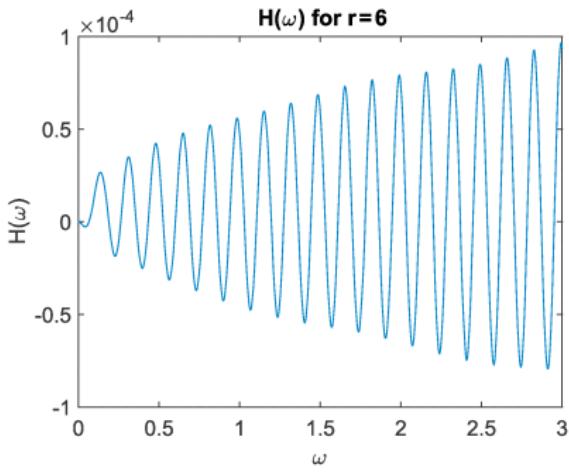


[Levi & Ori *PRD* **91** 104028 (2015)]

Generalized integrals

$$\mathcal{H}(\omega) = \int_{\omega=0}^{\omega} d\omega e^{i\omega\epsilon} \mathcal{G}_{\omega}(r)$$

$$\mathcal{H}_*(\omega) = \frac{1}{2} [\mathcal{H}(\omega) + \mathcal{H}(\omega + \nu/2)]$$



[Levi & Ori *PRD* **91** 104028 (2015)]

Reissner-Nordström black holes

$$f(r) = 1 - \frac{2M}{r} + \frac{Q^2}{r^2}$$

Reissner-Nordström black holes

$$f(r) = 1 - \frac{2M}{r} + \frac{Q^2}{r^2}$$

Horizons $f(r) = 0$

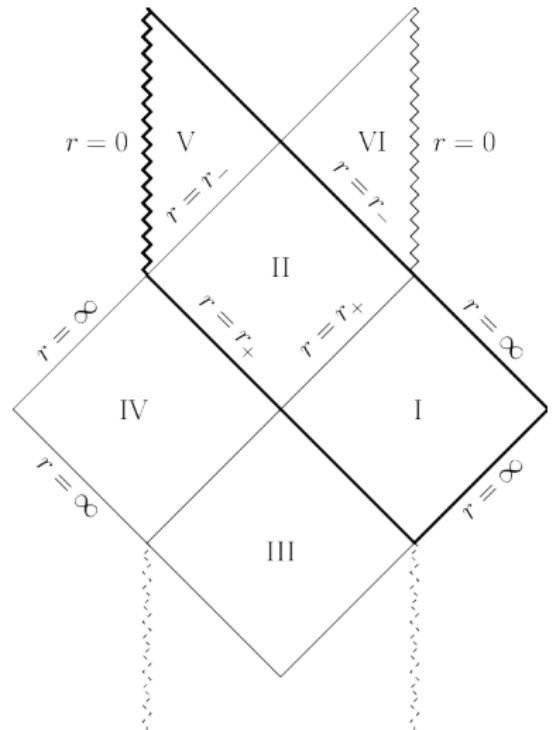
$$r_{\pm} = M \pm \sqrt{M^2 - Q^2}$$

Reissner-Nordström black holes

$$f(r) = 1 - \frac{2M}{r} + \frac{Q^2}{r^2}$$

Horizons $f(r) = 0$

$$r_{\pm} = M \pm \sqrt{M^2 - Q^2}$$



[Figure: Paston & Sheykin SIGMA **10** 003 (2014)]

Reissner-Nordström black holes

$$f(r) = 1 - \frac{2M}{r} + \frac{Q^2}{r^2}$$

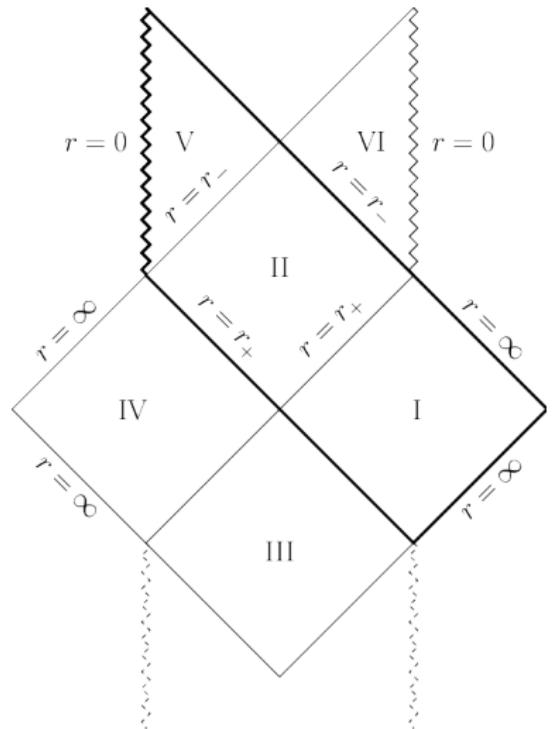
Horizons $f(r) = 0$

$$r_{\pm} = M \pm \sqrt{M^2 - Q^2}$$

Cosmic censorship

[Penrose (1974)]

- Cauchy horizon $r = r_-$
unphysical
- Classically unstable



[Figure: Paston & Sheykin SIGMA **10** 003 (2014)]

RSET at the Cauchy horizon in RN

[Zilberman, Levi & Ori *PRL* **124** 171302 (2020)

Zilberman & Ori *PRD* **104** 024066 (2021)]

RSET at the Cauchy horizon in RN

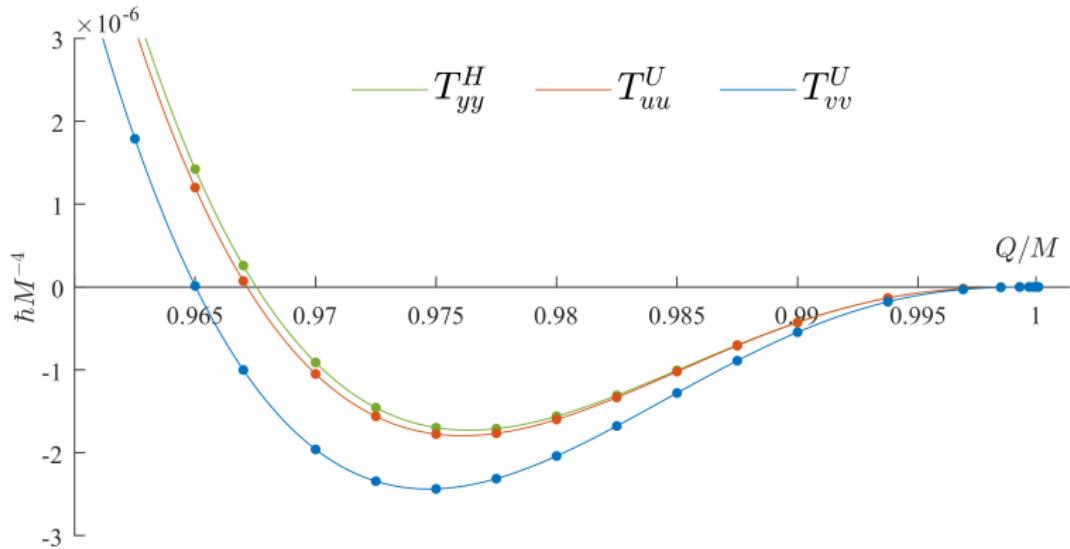
$$v = t + r_* \rightarrow \infty$$

[Zilberman, Levi & Ori *PRL* **124** 171302 (2020)

Zilberman & Ori *PRD* **104** 024066 (2021)]

RSET at the Cauchy horizon in RN

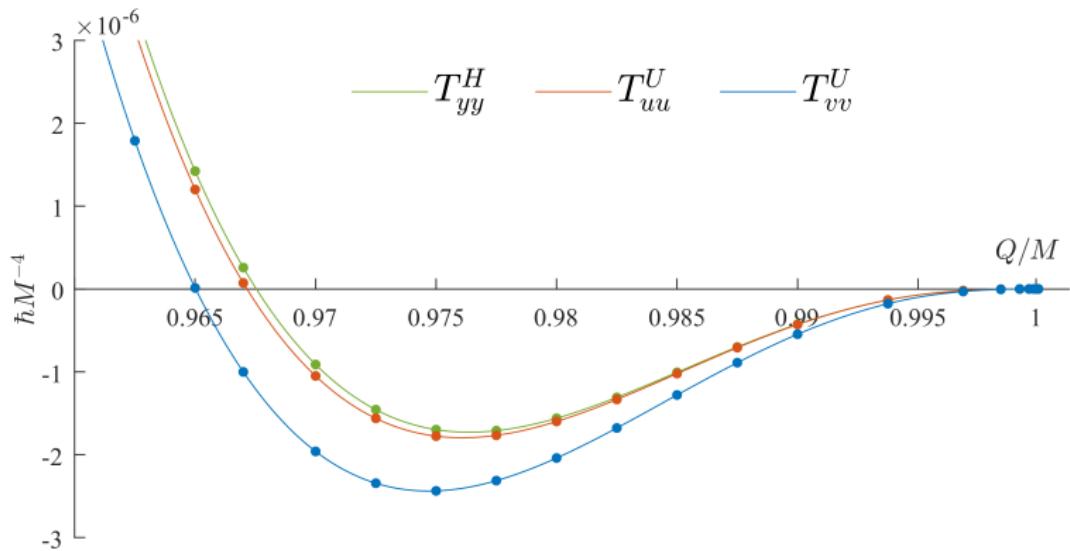
$$v = t + r_* \rightarrow \infty$$



[Zilberman, Levi & Ori *PRL* **124** 171302 (2020)
 Zilberman & Ori *PRD* **104** 024066 (2021)]

RSET at the Cauchy horizon in RN

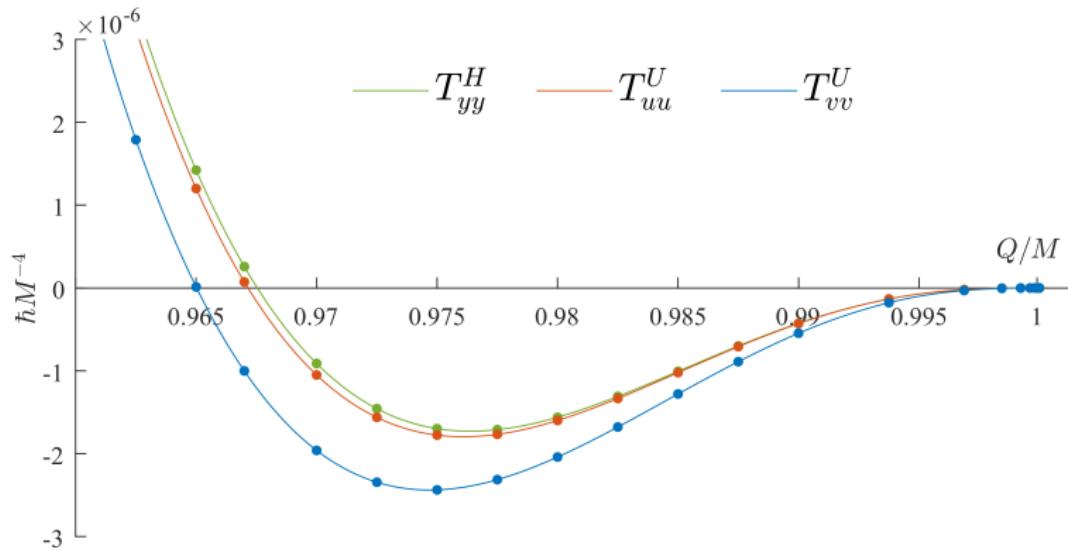
$$v = t + r_* \rightarrow \infty \quad V = -e^{-2\kappa_{r_*}} \rightarrow 0$$



[Zilberman, Levi & Ori *PRL* **124** 171302 (2020)
 Zilberman & Ori *PRD* **104** 024066 (2021)]

RSET at the Cauchy horizon in RN

$$v = t + r_* \rightarrow \infty \quad V = -e^{-2\kappa_r_*} \rightarrow 0 \quad \langle \hat{T}_{VV} \rangle \sim V^{-2} \langle \hat{T}_{vv} \rangle$$



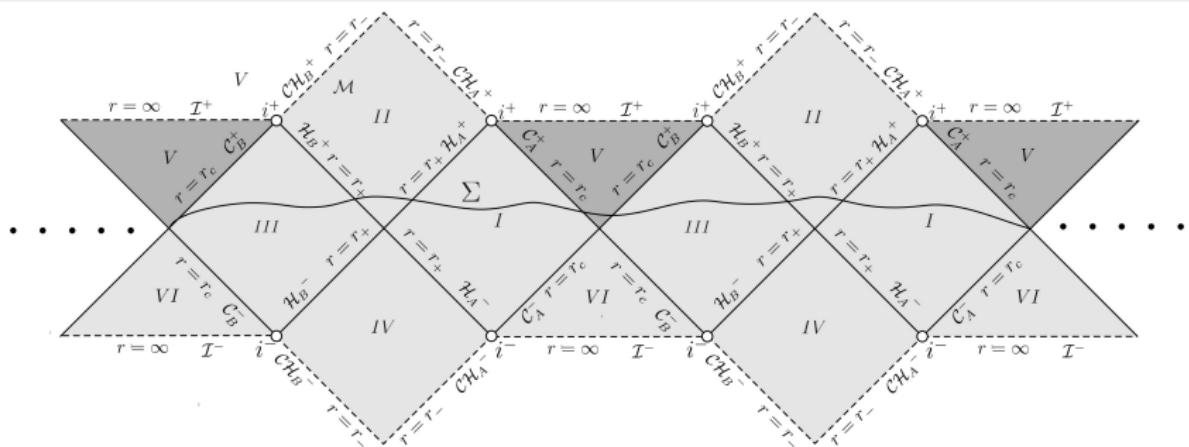
[Zilberman, Levi & Ori *PRL* **124** 171302 (2020)
 Zilberman & Ori *PRD* **104** 024066 (2021)]

Reissner-Nordström-de Sitter black hole

$$f(r) = 1 - \frac{2M}{r} + \frac{Q^2}{r^2} - \frac{\Lambda r^2}{3}$$

Reissner-Nordström-de Sitter black hole

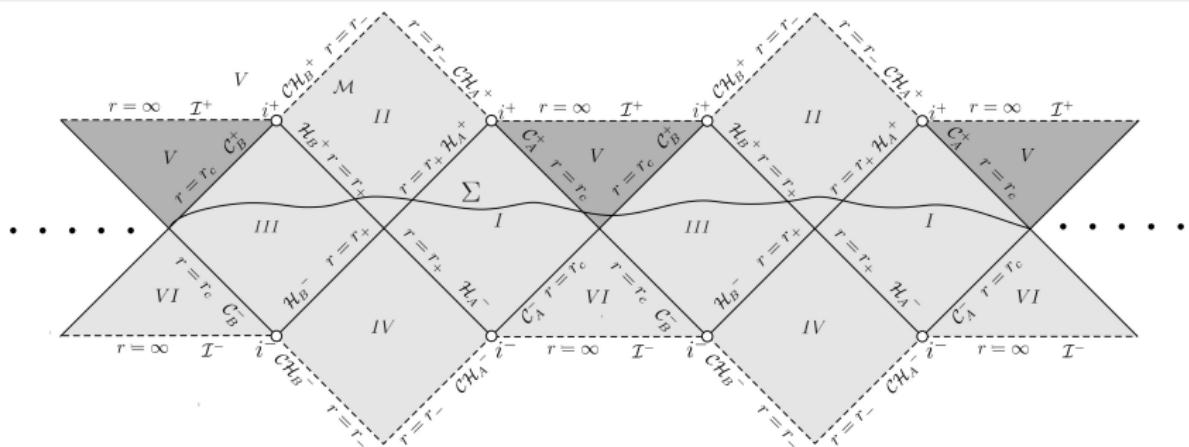
$$f(r) = 1 - \frac{2M}{r} + \frac{Q^2}{r^2} - \frac{\Lambda r^2}{3}$$



[Figure: Costa, Natario & Oliveira *Ann. H. Poincaré* **20** 3059 (2019)]

Reissner-Nordström-de Sitter black hole

$$f(r) = 1 - \frac{2M}{r} + \frac{Q^2}{r^2} - \frac{\Lambda r^2}{3}$$



[Figure: Costa, Natario & Oliveira *Ann. H. Poincaré* **20** 3059 (2019)]

Strong cosmic censorship fails classically for some RNdS black holes
 [Cardoso et al *PRL* **120** 031103 (2018)]

RSET at the Cauchy horizon in RNdS

RSET at the Cauchy horizon in RNdS

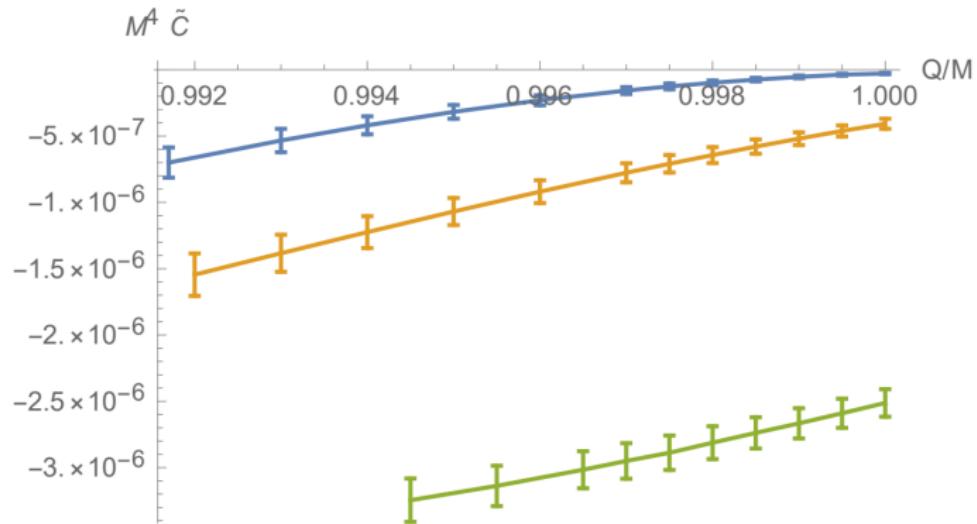
$$\langle \hat{T}_{VV} \rangle \sim \kappa_-^2 \tilde{C} |V|^{-2} + \dots \quad \text{as } V \rightarrow 0$$

[Hollands, Wald & Zahn *CQG* **37** 115009 (2020)]

RSET at the Cauchy horizon in RNdS

$$\langle \hat{T}_{VV} \rangle \sim \kappa_-^2 \tilde{C} |V|^{-2} + \dots \quad \text{as } V \rightarrow 0$$

[Hollands, Wald & Zahn *CQG* **37** 115009 (2020)]



[Hollands, Klein & Zahn *PRD* **102** 085004 (2020)]

Renormalized vacuum polarization and SET

Renormalized vacuum polarization and SET

Computing renormalized expectation values

- WKB-based method
- Extended coordinates method
- Pragmatic mode sum method

Renormalized vacuum polarization and SET

Computing renormalized expectation values

- WKB-based method
- Extended coordinates method
- Pragmatic mode sum method

Outlook

- Cosmic censorship

Renormalized vacuum polarization and SET

Computing renormalized expectation values

- WKB-based method
- Extended coordinates method
- Pragmatic mode sum method

Outlook

- Cosmic censorship
- Back-reaction problem

Semi-classical Einstein equations

$$G_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi \langle \hat{T}_{\mu\nu} \rangle$$